

Deformation of maximally supersymmetric Yang-Mills theory in dimensions 10. An algebraic approach.

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February 1, 2008

Abstract

We make a preliminary algebraic study of supersymmetric deformations of $N = 1$ Yang-Mills theory in dimension ten with an arbitrary gauge group. This is done in a context of Lie algebra deformation theory. The tangent space to the space of deformation is computed.

1 Introduction

In this paper we shall study the following problem. Suppose

$$\mathcal{L}(\nabla, \xi) = \langle F_{ij}, F_{ij} \rangle + \Gamma_{\alpha\beta}^i \langle \nabla_i \xi^\alpha, \xi^\beta \rangle \quad (1)$$

is a Lagrangian of supersymmetric Yang-Mills theory in dimension ten. Let θ_α be supersymmetry transformations.

*The work was partially supported by grants DE-FG02-90ER40542 and PHY99-0794

Some problem in physics can be formulated as a deformation problem: to find a certain deformation

$$\mathcal{L}_{\alpha'}(\nabla, \xi) = \mathcal{L}(\nabla, \xi) + \sum_{k \geq 1} \alpha'^k \mathcal{L}^k(\nabla, \xi) \quad (2)$$

where α' is a parameter of deformation. The deformed supersymmetry transformations

$$\theta_\alpha(\alpha') = \theta_\alpha + \sum_{k \geq 1} \theta_\alpha^k \alpha'^k \quad (3)$$

should leave the action corresponding to the density 2 invariant. The Lie algebra generated by $\theta_\alpha(\alpha')$, restricted to space of critical points of $\mathcal{L}_{\alpha'}$ is isomorphic to the supersymmetry algebra.

It is interesting to find (describe) all such deformations. We intend to split this problem into two sub problems:

A Find all infinitesimal deformations. By this we mean to solve the above problem up to order two in power series in α' . Thus we should look for $\mathcal{L}'(\nabla, \xi) = \mathcal{L}^1(\nabla, \xi)$ and $\theta'_\alpha = \theta_\alpha^1$ such that

$$\theta'_\alpha \mathcal{L}(\nabla, \xi) + \theta_\alpha \mathcal{L}(\nabla, \xi)' = \text{full derivative} \quad (4)$$

We may say that we are looking for on shell invariants of supersymmetry algebra.

We must also take into account a condition of on shell closure mod α'^2 of algebra of $\theta_\alpha(\alpha')$.

B Extend an infinitesimal deformation to an actual deformation.

In this paper we shall address sub problem **A**.

The gauge groups in our setup is $U(N)$, where N is large. Here is a list of the main results:

1 A system of equations (4) is overdetermined. It is not clear apriory that it has nonzero solutions at all. It seems possible however to give a formula for a general solution, which is $Spin(10)$ -invariant. It turns out that it is a sum of four terms

$$\mathcal{L}'(\nabla, \xi) = \mathcal{L}'_I(\nabla, \xi) + \mathcal{L}'_{II}(\nabla, \xi) + \mathcal{L}'_{III}(\nabla, \xi) + \mathcal{L}'_{IV}(\nabla, \xi) \quad (5)$$

$$\begin{aligned}
\mathcal{L}'_I(\nabla, \xi) = & \text{const}_I \text{tr} \left(\frac{1}{8} F_{mn} F_{nr} F_{rs} F_{sm} - \frac{1}{32} (F_{mn} F_{mn})^2 \right. \\
& + i \frac{1}{4} \xi^\alpha \Gamma_{m\alpha\beta} (\nabla_n \xi^\beta) F_{mr} F_{rn} \\
& - i \frac{1}{8} \xi^\alpha \Gamma_{mnr\alpha\beta} (\nabla_s \xi^\beta) F^{mn} F_{rs} \\
& + \frac{1}{8} \xi^\alpha \Gamma_{\alpha\beta}^m (\nabla_n \xi^\beta) \xi^\gamma \Gamma_{m\gamma\delta} (\nabla_n \xi^\delta) \\
& \left. - \frac{1}{4} \xi^\alpha \Gamma_{\alpha\beta}^m (\nabla_n \xi^\beta) \xi^\gamma \Gamma_{n\gamma\delta} (\nabla_m \xi^\delta) \right) \\
& \left(\text{The formula for } \mathcal{L}'_I(\nabla, \xi) \text{ was borrowed from [1]} \right) \\
\mathcal{L}'_{II}(\nabla, \xi) = & \text{const}_{II} \mathcal{L}_3(\nabla, \xi) \\
& \left(\text{formula 3.1 page 6 in [6] is too long to be presented here} \right)
\end{aligned} \tag{6}$$

It is α'^2 coefficient of power series (2.63) in [1] .

At the moment we have only partial information about $\mathcal{L}'_{III}(\nabla, \xi)$ that it one of two **susy** invariants of degree 5 in α' . See section (4.2) however.

Suppose $n(\nabla, \xi)$ is a noncommutative polynomial in covariant derivatives of curvature and spinor field. Define $\nu(\nabla, \xi) = \text{tr} n(\nabla, \xi)$.

Define

$$\mathcal{L}'_{IV}(\nabla, \xi) = \epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}} \nu(\nabla, \xi) \tag{7}$$

where ϵ is the Levi-Chevita tensor.

The theories with smaller supersymmetries which allow off shell formulation with manifest supersymmetries admit a simple construction of deformation $\mathcal{L}(\nabla, \xi) \rightarrow \mathcal{L}(\nabla, \xi) + \alpha' \mathcal{L}'(\nabla, \xi)$ of supersymmetric Lagrangian. $\mathcal{L}'(\nabla, \xi)$ can be chosen as an arbitrary superfunction in chiral superfields. In this case ν becomes one of the components of the chiral superfunction. The operator $\epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}}$ can be interpreted as "chiral odd integration" (see [8] chapter 10). The analogy with theories with low number of supersymmetries is not accidental. In the course of the proof we used pure spinor approach of

Howe-Berkovits (its homological version), which could be considered as the best approximation to manifestly susy-covariant formulation of Yang-Mills theory .

2 Collection θ'_α uniquely determine $\mathcal{L}'(\nabla, \xi)$ and visa versa.

Since everything is covariant with respect to the group of translations we may assume that the coefficients of the Lagrangian \mathcal{L}' are some constants. We can associate a degree with every Lagrangian $\mathcal{L}' = \text{trn}(\nabla, \xi)$ by the rule $\deg \nabla_i = 2, \deg \xi^\alpha = 3$. The degree of all constants is zero. As a result $\deg \mathcal{L}(\nabla, \xi) = 8$ for $\mathcal{L}(\nabla, \xi)$ defined in 1. The relation of the degree \deg with more conventional degree in α' which we denote by $\deg_{\alpha'}$ is

$$\deg_{\alpha'} = \frac{\deg - 8}{4} \quad (8)$$

3 One can form a generating function $\tilde{a}(t) = \sum_{k \geq 0} \tilde{a}_k t^k$. The coefficients a_k are dimensions of linear spaces spanned by infinitesimal Lagrangians \mathcal{L}' of degree $\deg = k$, which satisfy (4), defined up to a field redefinition. Then

$$\tilde{a}(t) = t^8(a(t) - 126t^{-2} - 144t) \quad (9)$$

The formula for $a(t)$ is given in (59).

A typical Lagrangian \mathcal{L}' (no condition (4) is necessary at this point) is defined for a gauge group $U(N)$, because it involves products (of derivatives) of curvature, like in $\text{tr}(F_{ij}F_{jk}F_{kl}F_{li})$ in (6). The last Lagrangian is not defined for exceptional groups because it requires an additional data - a representation in $U(N)$. A possible restriction on \mathcal{L}' is that it is defined for any Lie algebra, i.e. it is of the form

$$\sum_r (m_{r,1}(\nabla, \xi), m_{r,2}(\nabla, \xi)) \quad (10)$$

In the last formula $(., .)$ is the invariant dot product on the Lie algebra of the gauge group \mathfrak{g} , $m_{r,1}, m_{r,2}$ are commutators in covariant derivatives of curvature and spinor fields. We call Lagrangians written in (10) the Lagrangians of the Lie type. In this setup we can form a generating function $\tilde{l}(t)$ in analogy with $\tilde{a}(t)$. Then

$$\tilde{l}(t) = t^8(l(t) - 144t) \quad (11)$$

The formula for $l(t)$ is given in (59).

4 One can define a series of types of Lagrangians that are similar to the Lie type. They can be defined by the formula

$$\mathcal{L}'(\nabla, \xi) = \sum_{\sigma \in S_p} \sum_r \pm \text{tr}(m_{r, \sigma(1)}(\nabla, \xi) \dots m_{r, \sigma(p)}(\nabla, \xi)) \quad (12)$$

The commutators $m_{r,i}(\nabla, \xi)$ $1 \leq i \leq p$ are defined as before. The auxiliary degree of such Lagrangian is equal to p . One can form a generating function $l(t, u) = \sum_{k,p} l_{kp} t^k u^p$, with $l_{kp} = \dim L_{kp}$ -dimensions of space of Lagrangians of bidegree k, p up to field redefinition and satisfying (4). We do not provide a formula for $l(t, u)$ (our technique allows to do that but the formula becomes messy). Instead we tabulated below the first few coefficients of a related function $\tilde{l}^{Spin(10)}(t, u)$. The coefficient of $l^{Spin(10)}(t, u)$ are dimension of $Spin(10)$ -invariants in spaces L_{kp} . In fact it is convenient to use $deg_{\alpha'}$ instead of deg at this point, because

$$L_{k,p}^{Spin(10)} = 0 \quad k \not\equiv 0 \pmod{4}$$

. A transition $deg \rightarrow deg_{\alpha'}$ in the degrees of generating function manifests in a change of variables $l^{Spin(10)}(t^{\frac{1}{4}}, u)t^{-2} = \tilde{l}^{Spin(10)}(t, u)$ (t^{-2} -factor is for agreement with (8)). The coefficients \tilde{l}_{kp} of $\tilde{l}^{Spin(10)}(t, u)$ are

	1	2	3	4	5	6	7	8	$k = deg_{\alpha'}$
2			1		1	3	18	172	...
3							13	281	...
4		1			1	2	20	267	...
5							1	68	...
6							1	17	...
7									...
p

(13)

Let us say a few words about one important feature of Yang-Mills theory which was not accommodated in the present framework.

The Lagrangian 1 admits additional "trivial" translational supersymmetries: $\tilde{\theta}_\alpha \nabla_i = 0$, $\tilde{\theta}_\alpha \chi^\beta = \delta_\alpha^\beta$. A systematic treatment of these we defer to future publications.

The methods of this paper are similar to [15],[16], [13]. The principal ingredient of this paper is a Lie algebra YM . One can think about it as of a universal solution of D=10, N=1 Yang-Mills equations.

It comes about as follows: we replace fields by generators, equations of motion by relations. In case of Yang-Mills equations this way we get YM algebra. The supersymmetry generators also can be lifted to some formal variables. They act on YM mimicking formulas of component formalism (21). We denote by L the Lie algebra of derivations they generate. This algebra was introduced in [15] and studied in [13]. It turns out that L contains YM as an ideal. There is a projection $L \rightarrow \mathfrak{su}(2,2)$ with a kernel TYM . It turns out that there is a tower of inclusions $TYM \subset YM \subset L$.

Suppose that we constructed some supersymmetric deformation of D=10, N=1 Yang-Mills theory. The outlined above procedure will provide us with deformations $YM_{\alpha'}$ and $L_{\alpha'}$. Our main assumption is that in process of deformation dimensions of invariantly defined space do not jump. In particular a sequence of ideals $TYM_{\alpha'} \subset YM_{\alpha'} \subset L_{\alpha'}$ survives in the process of deformation.

Such algebraization allows us to use powerful methods of deformation theory of algebraic systems (Lie and associative algebras).

The paper is organized as follows: We collected all necessary notations and definition in sections (1.1), 1.2.

In section (2) we setup a stage. We define precisely how we understand a supersymmetric deformation of Yang-Mills theory. We have several flavors of deformations that are classified by types. We distinguish two types: associative (**A**) and Lie (**L**) types. We also introduce an independent **Lg** condition - the deformed Yang-Mills theory should have a Lagrangian. We also introduce deformation complexes.

In section (3) we make the main reduction in deformation complexes, replacing the standard one by much smaller. It enables us to do the computations. We use pure spinors in essential way. The most important technical proposition in this section is (20).

In section 4 we give partial justification of the claim **1** from the introduction

In section (5) we justify claims **3,4**.

In section (6) we verify claim **2** and finish with **1**.

We also provide an extensive appendix where some algebraic facts are collected and proofs of some technical statements are outlined.

In particular, we have a section 7.1 on quadratic algebras where the reader can find explanations of the main reduction from section (3).

We also have a section (7.4) on equivariant cyclic homology. Some facts collected in it are useful for better understanding material from section 4, (5).

Acknowledgment 1 *The author would like to thank MPI, KITP, IAS, where the most of the work has been done. He also would like to thank N.Nekrasov, M. Rocek, A.S.Schwarz, D.Sullivan, for useful discussions and comments . Also I would like to thank B. Lowson and M.Roczek for opportunity to present some of this material at "2005 Simons workshop".*

1.1 Notations

All linear spaces in this note are defined over complex numbers \mathbb{C} .

An abbreviation for a complex

$$\cdots \rightarrow A_i \rightarrow A_{i+1} \rightarrow \cdots$$

is A^\bullet . There is a standard shift operation on complexes $(A^\bullet[n])_i = (A^\bullet)_{i+n}$, $d_{A^\bullet[n]} = (-1)^n d_{A^\bullet}$.

Let C be an algebra, $\varepsilon : C \rightarrow \mathbb{C}$ a homomorphism (augmentation). Denote $IC = \text{Ker} \varepsilon$

We denote $\text{Sym}(W) = \bigoplus_{i \geq 0} \text{Sym}^i(W)$ a symmetric algebra of a linear space W . Denote $a \bullet b$ a product of elements $a, b \in \text{Sym}(W)$. The object $\Lambda(W) = \bigoplus_{i \geq 0} \Lambda^i(W)$ is the Grassman algebra on W . If W is a graded vector space that $\text{Sym}(W), \Lambda(W)$ are defined conforming to the sign rules.

We denote a bracket in a Lie algebra by $[\cdot, \cdot]$ or $\{\cdot, \cdot\}$. We shall use uniform notations for commutators and anticommutators.

We use Einstein convention of summation over repeated indices and do not use \sum sign. In case when \sum is present we perform a summation over non-tensor indices, the range of summation can be guessed from the context.

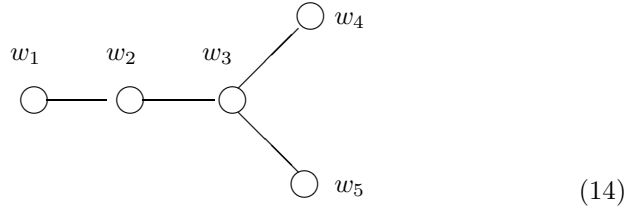
Let V be 10-dimensional linear space over complex numbers. It is equipped with a symmetric nondegenerate dot-product invariant with respect to $Spin(10)$. Let S be an irreducible complex spinor representation of $Spin(10)$. Vectors v_1, \dots, v_{10} is an orthonormal basis of V , χ^α $\alpha = 1, \dots, 16$ is a basis of S . θ_α is a dual basis of S^* .

This data allows us define Γ -matrices $\Gamma_{\alpha\beta}^i, \Gamma_i^{\alpha\beta}, \Gamma_{\alpha}^{\beta ij}, \Gamma_{\alpha\beta}^{i_1 i_2 i_3}, \dots, \Gamma_{i_1 i_2 i_3 i_4 i_5}^{\alpha\beta} \dots$. See [13] for discussion of spinors and Γ -matrices.

All tensors involved in our considerations are build from vector representation V (latin vector indices), irreducible spinor representations S, S^* (Greek spinor indices) of $Spin(10)$. A presence of a dot-product permits us to make no distinction between lower and upper vector indices. We however shall make a careful distinction between lower and upper spinor indices.

Denote $f \circ g$ a composition of maps: $(f \circ g)(x) = f(g(x))$.

We use the standard convention labeling representation of a semisimple group by it highest weight. The Dynkin graph of the group $Spin(10)$ is



The labels w_i correspond to coordinates of the highest weight. We encode a representation that is labeled by Dynkin diagram with labels as above by an array $[w_1, w_2, w_3, w_4, w_5]$. Our convention is that spinor representation $S_1 \subset S$ is equal to irreducible representation with highest weight $[0, 0, 0, 0, 1]$, the tautological representation in \mathbb{C}^{10} is equal to $[1, 0, 0, 0, 0]$, the adjoint representation is equal to $[0, 1, 0, 0, 0]$... A general representation is then $\bigoplus_{w_i \geq 0} a_{w_1, \dots, w_5} [w_1, \dots, w_5]$, $a_{w_1, \dots, w_5} \in \mathbb{Z}_{\geq 0}$ are the multiplicities.

1.2 Main definitions.

Definition 2 *YM algebra is a quotient of a free graded Lie algebra $\widetilde{YM} = \text{Free} \langle v_1, \dots, v_n, \chi^1, \dots, \chi^{16} \rangle$, $\deg(v_i) = 2, \deg(\chi^\alpha) = 3$ by an ideal. The ideal is generated by relations*

$$\tilde{v}_m = [v_s, [v_s, v_m]] - \frac{1}{2} \Gamma_{\alpha\beta}^m [\chi^\alpha, \chi^\beta] \quad (15)$$

$$\tilde{\chi}_\alpha = \Gamma_{\alpha\beta}^s [v_s, \chi^\beta] \quad (16)$$

Definition 3 *A Lie algebra L is a quotient of a free graded Lie algebra $\text{Free} \langle \theta_1, \dots, \theta_{16} \rangle$. $\deg \theta_i = 1$. The generators of the ideal of relations are*

$$\Gamma_{i_1 \dots i_5}^{\alpha\beta} [\theta_\alpha, \theta_\beta] = 0 \quad (17)$$

From [15], [16] we know that YM is a subalgebra of Lie algebra L . The embedding $\rho : YM \rightarrow L$ is defined by the formulas

$$\begin{aligned} \rho(v_i) &= \Gamma_i^{\alpha\beta} [\theta_\alpha, \theta_\beta] \\ \rho(\chi^\alpha) &= \Gamma^{\alpha\beta s} [\rho(v_s), \theta_\beta] \end{aligned} \quad (18)$$

Denote

$$L = \bigoplus_{i \geq 1} L_i \quad (19)$$

-decomposition into graded pieces. It was proved in [13] that

$$YM = \bigoplus_{i \geq 2} L_i \quad (20)$$

In the following we identify YM with subalgebra of L such that the following formula holds $[\theta_\alpha, \theta_\beta] = \Gamma_{\alpha\beta}^i v_i$

The algebra YM is an ideal in L . We interpret the adjoint action of $\theta_\alpha \in L_1$ on YM as an action of supersymmetries:

$$\begin{aligned} \theta_\alpha v_i &= \Gamma_{\alpha\beta i} \chi^\beta \\ \theta_\beta \chi^\alpha &= \Gamma_\beta^{\alpha ij} F_{ij} \end{aligned} \quad (21)$$

Following [16] and [13] denote

$$TYM = \bigoplus_{i \geq 3} L_i \quad (22)$$

From [16] we know that TYM is generated as a Lie algebra by elements of the form

$$\begin{aligned} [v_{i_1}, \dots, [v_{i_k}, F_{ij}] \dots] \quad F_{ij} = [v_i, v_j] \\ [v_{i_1}, \dots, [v_{i_k}, \chi^\alpha] \dots] \end{aligned} \quad (23)$$

Definition 4 Denote by $\mathfrak{su}\mathfrak{sh}$ a graded variant of Lie algebra of supersymmetries in dimension ten. It is \mathbb{Z} -graded algebra $\mathfrak{su}\mathfrak{sh}_1 = S^* = \langle \theta_1, \dots, \theta_{16} \rangle$, $\mathfrak{su}\mathfrak{sh}_2 = V = \langle v_1, \dots, v_{10} \rangle$. After reduction of grading modulo two we get a standard \mathbb{Z}_2 grading on $\mathfrak{su}\mathfrak{sh}$. The space $\mathfrak{su}\mathfrak{sh}_2$ is the center. The commutator of $\theta_\alpha \in \mathfrak{su}\mathfrak{sh}_1$ is defined by the formula $[\theta_\alpha, \theta_\beta] = 2\Gamma_{\alpha\beta}^i v_i$.

Let us remind the standard relation of YM -algebra to the classical Yang-Mills theory.

Let ∇ be a connection in a principle $U(N)$ -bundle on \mathbb{R}^{10} . We assume that a choice of coordinates x_i on \mathbb{R}^{10} is given in which a metric has a diagonal form dx_i^2 . Denote $S^* \otimes \mathfrak{u}(N)$ a tensor product of irreducible complex dual spinor and adjoint bundles. Let $\tilde{\xi}_\alpha$, $\alpha = 1, \dots, 16$ be a set of sections which form a basis of S (pointwise). Then any section of $S \otimes \mathfrak{u}(N)$ can be presented as $\tilde{\xi}_\alpha \xi^\alpha$, where ξ^α is sixteen $\mathfrak{u}(N)$ -valued functions. Denote $\nabla_i = \nabla_{\frac{\partial}{\partial x_i}}$. Let us assume that the Γ -matrices $\Gamma_{\alpha\beta}^i$ in the bases are translationary invariant and equal to the standard Γ -matrices used in definition (2).

Connection in the principal bundle define covariant derivative in any associated bundle, which we denote by the same letter ∇ . We shall be interested in $\Lambda(\Theta) \otimes T$, where T is the fundamental representation of $U(N)$. We can interpret $\Lambda(\Theta)$ as function on auxiliary odd space of parameters Θ^* . As it is common in supergeometry connection ∇ has coefficients in $\Lambda(\Theta)$. The same applies to spinor: ξ^α is a matrix spinor with coefficients in $\Lambda(\Theta)$, i.e. section of $S^* \otimes \Lambda(\Theta) \otimes \mathfrak{u}(N)$.

Levi-Chevita connection, together with connection ∇ define covariant derivatives ∇_i acting in sections of $\Lambda(\Theta) \otimes T$, which we can consider as differential operators of the first order. Any odd section of $\Lambda(\Theta) \otimes \mathfrak{u}(N)$ defines a matrix

operator with $\Lambda(\Theta)$ entries acting on $\Lambda(\Theta) \otimes T$. For a given choice of connection ∇_i and spinors ξ^α they generate a subalgebra in a (super)Lie algebra of differential operators acting in $\Lambda(\Theta) \otimes T$.

An assignment

$$\begin{aligned} v_i &\rightarrow \nabla_i \\ \chi^\alpha &\rightarrow \xi^\alpha \end{aligned} \tag{24}$$

is a homomorphism of YM -algebra to the algebra of differential operators if and only if ∇_i, ξ^α is a solution of classical Yang-Mills equation

$$\begin{aligned} \nabla_i F_{ij} &= \frac{1}{2} \Gamma_{\alpha\beta}^j [\xi^\alpha, \xi^\beta], \quad [\nabla_i, \nabla_j] = F_{ij} \\ \Gamma_{\alpha\beta}^i \nabla_i \xi^\beta &= 0 \end{aligned} \tag{25}$$

Standard supersymmetry transformations can be obtained from (21) after a substitution (24). These transformations satisfy relation (17) if ∇_i, ξ^α is a solution of classical Yang-Mills equation.

2 Definition of supersymmetric deformation of YM algebra

Fix Lie algebras \mathfrak{n} and \mathfrak{l} .

Definition 5 *A Lie algebra \mathfrak{g} belongs to the class $\mathbf{L}(\mathfrak{n}, \mathfrak{l})$ if \mathfrak{g} is an extension of \mathfrak{n} by \mathfrak{l} , i.e. fits into exact sequence*

$$0 \rightarrow \mathfrak{l} \rightarrow \mathfrak{g} \rightarrow \mathfrak{n} \rightarrow 0$$

.

By construction we have $\mathfrak{su}\mathfrak{s}\mathfrak{h} = L/TYM$. It means that $L \in \mathbf{L}(\mathfrak{su}\mathfrak{s}\mathfrak{h}, TYM)$

In the following text α' is a formal parameter of deformation. All maps involved in the construction are $\mathbb{C}[[\alpha']]$ -linear. All deformations are flat modules over $\mathbb{C}[[\alpha']]$. We have an isomorphism of $\mathbb{C}[[\alpha']]$ modules $L_{\alpha'} = L \hat{\otimes} \mathbb{C}[[\alpha']]$. Lie

algebra has a topology defined by filtration $F^i = \bigoplus_{k \geq i} L_k$ (the reader can consult [16] about completions).

To shorten the notations we denote a trivial deformation of Lie algebra \mathfrak{g} as $\mathfrak{g}(\alpha')$. It is isomorphic to $\mathfrak{g} \hat{\otimes} \mathbb{C}[[\alpha']]$ not just as a $\mathbb{C}[[\alpha']]$ -module but as a Lie algebra.

Definition 6 *We shall be interested in deformations $L_{\alpha'}$ of L which contains a Lie subalgebra $TYM_{\alpha'}$ such that there is a short exact sequence of algebras*

$$0 \rightarrow TYM_{\alpha'} \rightarrow L_{\alpha'} \xrightarrow{p} \mathfrak{su}\mathfrak{h}(\alpha') \rightarrow 0 \quad (26)$$

and $\mathfrak{su}\mathfrak{h}$ stays undeformed. Two deformations $L_{\alpha'}$ and $L'_{\alpha'}$ are equivalent if there is an isomorphism $\eta : L_{\alpha'} \rightarrow L'_{\alpha'}$, which map $TYM_{\alpha'}$ into $TYM_{\alpha'}$. We shall call such class of deformations- $\mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM_{\alpha'})$ -deformations.

Remark 7 *The algebra TYM is free (see [16]), therefore are rigid. The class $\mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM_{\alpha'})$ coincides with $\mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$.*

Definition 8 $\mathbf{L}^{Spin(10)}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$ *a subclass of $\mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$ of $Spin(10)$ -equivariant deformation. The $Spin(10)$ representation content of L_{α} must coincide with the content of L .*

A deformation of L of type $\mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$ gives rise to a deformation of YM algebra. Indeed we can take $YM_{\alpha'} = [L_{\alpha'}, L_{\alpha'}] = \{\sum_i [a_i, b_i] | a_i, b_i \in L_{\alpha'}\}$. $YM_{\alpha'}$ is an ideal in $L_{\alpha'}$. We recover the action of supersymmetries from adjoint action of generators of $L_{\alpha'}$ on $YM_{\alpha'}$.

An algebra $L_{\alpha'} \in \mathbf{L}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$ admits a decreasing filtration $F^i = [L_{\alpha'}, \dots, [L_{\alpha'}, L_{\alpha'}] \dots] - i$ times. Except few possible pathological cases $\bigcap_i F^i = 0$, the last condition holds true if $L_{\alpha'} \in \mathbf{L}^{Spin(10)}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$.

The formal definition of deformation can be translated into the language of connections. The number of generators and relations of $YM_{\alpha'}$ is the same as of YM , if $L_{\alpha'} \in \mathbf{L}^{Spin(10)}(\mathfrak{su}\mathfrak{h}(\alpha'), TYM(\alpha'))$. This is because of the mentioned filtration. On YM the filtration is equal to $F^i = \bigoplus_{k \geq i} YM_k$. We identify $YM_{\alpha'}$ with $YM \otimes \mathbb{C}[[\alpha']]$ as $\mathbb{C}[[\alpha']]$ -modules via identity transformations. This

enables us to interpret algebraic generators of YM as generators of $YM_{\alpha'}$ over $\mathbb{C}[[\alpha']]$. The deformed relations must then be of the form

$$\begin{aligned} [v_i[v_i, v_j]] - \frac{1}{2}\Gamma_{\alpha\beta}^j[\chi^\alpha, \chi^\beta] &= \alpha' r_j(\alpha') \\ \Gamma_{\alpha\beta}^i[v_i, \chi^\beta] &= \alpha' s_\alpha(\alpha') \end{aligned} \quad (27)$$

The elements $r_j(\alpha'), s_\alpha(\alpha')$ are some formal power series with coefficients in ideal of free Lie algebra \widetilde{YM} generated by (23). After substitution (24) the remainders $r_j(\alpha'), s_\alpha(\alpha')$ will become formal powers series with coefficients in Lie algebra, generated by elements (28)

$$\begin{aligned} \nabla_{i_1} \dots \nabla_{i_k} F_{st} \\ \nabla_{i_1} \dots \nabla_{i_k} \xi^\alpha \end{aligned} \quad (28)$$

for $k \geq 0$.

We say that ∇_i, ξ^α is a solution of the deformed Yang-Mills equation if ∇_i, ξ^α to satisfy (27) after substitution (24).

There is a standard language that was proved to be useful in solving deformation problems of above type. It is the language of deformation theory of Lie algebras

Definition 9 Suppose \mathfrak{g} is an arbitrary Lie algebra and N is a \mathfrak{g} -module. It is a homomorphism $\rho : \mathfrak{g} \rightarrow \text{End}(N)$. There is a complex

$$C^k(\mathfrak{g}, N) = \text{Hom}(\Lambda^k(\mathfrak{g}), N) \quad (29)$$

, called Cartan-Chevalley complex. The differential $d : C^k(\mathfrak{g}, N) \rightarrow C^{k+1}(\mathfrak{g}, N)$ is defined by the formula:

$$\begin{aligned} (dc)(l_1, \dots, l_{k+1}) &= \sum_{i=1}^{k+1} (-1)^i \rho(l_i) c(l_1, \dots, \hat{l}_i, \dots, l_{k+1}) + \\ &+ \sum_{i < j} (-1)^{i+j-1} c([l_i, l_j], l_1, \dots, \hat{l}_i, \dots, \hat{l}_j, \dots, l_{k+1}) \end{aligned} \quad (30)$$

The cohomology of this complex is denoted by $H^k(\mathfrak{g}, N)$.

A complex

$$C_k(\mathfrak{g}, N) = \Lambda^k(\mathfrak{g}) \otimes N \quad (31)$$

has a differential

$$\begin{aligned}
d : C_k(\mathfrak{g}, N) &\rightarrow C_{k-1}(\mathfrak{g}, N) \\
d(n \otimes l_1 \wedge \cdots \wedge l_k) &= \sum_{i=1}^k (-1)^i \rho(l_i) n \otimes l_1 \wedge \cdots \wedge \hat{l}_i \wedge \cdots \wedge l_k + \\
&+ \sum_{i < j} (-1)^{i+j-1} n \otimes [l_i, l_j] \wedge \cdots \wedge \hat{l}_i \wedge \cdots \wedge \hat{l}_j \wedge \cdots \wedge l_k \quad l_s \in \mathfrak{g}, n \in N
\end{aligned} \tag{32}$$

The cohomology of this complex is denoted by $H_k(\mathfrak{g}, N)$. There is an obvious extension of these constructions to \mathbb{Z}_2 or \mathbb{Z} graded category. For details the reader can consult [9].

For deformation purposes we shall be interested in adjoint representation $N = \mathfrak{g}$ or adjoint representation in universal enveloping $U(\mathfrak{g})$. We shall give some illustrations of usefulness of complex $C^\bullet(\mathfrak{g}, N)$. For simplicity in the introductory discussion we assume that \mathfrak{g} is a purely even Lie algebra.

Let us start with $c \in C^1(\mathfrak{g}, \mathfrak{g})$. It is easy to see that the condition $dc = 0$ is the equation $c([l_1, l_2]) = [l_1, c(l_2)] + [c(l_1), l_2]$. It is a condition that c is a derivation of \mathfrak{g} : $c \in \text{Der}(\mathfrak{g})$. There is a class of trivial derivations $\text{In}(\mathfrak{g})$ -so called inner derivations $c_a(l) = [a, l]$. It is natural to work with a quotient

$$\text{Der}(\mathfrak{g})/\text{In}(\mathfrak{g}) = \text{Out}(\mathfrak{g}) \tag{33}$$

The later linear space by definition coincides with $H^1(\mathfrak{g}, \mathfrak{g})$.

As an exercise the reader can check that $H^0(\mathfrak{g}, \mathfrak{g})$ coincides with the center of \mathfrak{g} .

Any derivation of \mathfrak{g} defines a derivation of universal enveloping $U(\mathfrak{g})$. The converse is not true. If one would like to understand derivations of $U(\mathfrak{g})$, one has to replace N by $U(\mathfrak{g})$ in (29) and compute the first cohomology.

The groups $H^k(\mathfrak{g}, \mathfrak{g})$ are a direct summands in $H^k(\mathfrak{g}, U(\mathfrak{g}))$.

The linear space $H^2(\mathfrak{g}, \mathfrak{g})$ can be interpreted as a space of nonequivalent infinitesimal deformations of \mathfrak{g} . Indeed if we expand a deformed bracket $[\cdot, \cdot]_{\alpha'}$ into formal series in α' , we shall get

$$[\cdot, \cdot]_{\alpha'} = [\cdot, \cdot] + \alpha' \gamma_1(\cdot, \cdot) + \alpha'^2 \gamma_2(\cdot, \cdot) + \dots \tag{34}$$

Denote $\gamma(.,.) = \gamma_1(.,.)$. The map $a_1 \otimes a_2 \otimes a_3 \rightarrow \sum_{\sigma \in \mathbb{Z}_3 \subset S_3} [[a_{\sigma(1)}, a_{\sigma(2)}]_{\alpha'}, a_{\sigma(3)}]_{\alpha'}$ (which is zero for a Lie algebra) has its first Taylor coefficient equal to

$$\sum_{\sigma \in \mathbb{Z}_3 \subset S_3} [\gamma(a_{\sigma(1)}, a_{\sigma(2)}), a_{\sigma(3)}] + \gamma([a_{\sigma(1)}, a_{\sigma(2)}], a_{\sigma(3)}) \quad (35)$$

It is equal to zero precisely when γ is two-cocycle. Denote a space of two cocycles by Z^2 . There is a subspace $B^2 \subset C^2(\mathfrak{g}, \mathfrak{g})$, which is generated by infinitesimal action of Lie algebra of coordinate change $\mathfrak{gl}(\mathfrak{g})$ on the bracket, viewed as an element $e \in C^2(\mathfrak{g}, \mathfrak{g})$. We may identify $\mathfrak{gl}(\mathfrak{g})$ with $C^1(\mathfrak{g}, \mathfrak{g})$. Denote the action of $l \in \mathfrak{gl}(\mathfrak{g})$ on e by le . Then $le = d(l)$. We see that $dB^2 = 0$. As a result a tangent space to the space of deformations of \mathfrak{g} is equal to $Z^2/B^2 = H^2(\mathfrak{g}, \mathfrak{g})$. Recalling discussion of derivations the reader should not be surprised to know that the space $H^2(\mathfrak{g}, U(\mathfrak{g}))$ classifies infinitesimal deformations of $U(\mathfrak{g})$.

Deformation of the bracket $[a, b]_{\alpha'} = [a, b] + \alpha' \gamma(a, b) + \dots$ of L which satisfies (26) must satisfy $p\gamma = 0$. Hence $\text{Im} \gamma \subset \text{TYM}$. The following proposition becomes obvious

Proposition 10 *Infinitesimal deformations of L which can be put into short exact sequence (26) are classified by elements of $H^2(L, \text{TYM})$, infinitesimal deformations $\mathbf{L}^{Spin(10)}(\mathfrak{su}\mathfrak{sh}(\alpha'), \text{TYM}(\alpha'))$ are classified $H^2(L, \text{TYM})^{Spin(10)}$. In both cases, the degrees of deformations are even.*

The elements of $H^i(L, \text{TYM})$ are graded. The grading of TYM starts with three. For a homogeneous cocycle $\gamma(a, b) \in C^2(L, \text{TYM})$ its degree can be determined by a formula $\deg(\gamma(\theta_\alpha, \theta_\beta)) - 2$. It is greater than zero. An additional restriction on the cocycle is that its degree must be even.

It is also useful to study cohomology of $H^2(L, U(\text{TYM}))$. We provide below a somewhat cumbersome definition of deformation of algebra $U(L)$ which leads to the space of infinitesimal deformations equal to $H^2(L, U(\text{TYM}))$. A difficulty is that $U(L)_{\alpha'}$ is no longer a universal enveloping. It however retains certain properties of universal enveloping, which enables us to interpret the action of $\theta_\alpha(\alpha')$ by commutators on $U(\text{TYM})$ as supersymmetries.

One can define a Lie algebra $Out(C)$ for an associative algebra C . The definition mimics to the Lie algebra case (33).

The linear space $L_1 + L_2 \subset L$ normalizes Lie subalgebra TYM . It also normalizes $U(TYM) \subset U(L)$. This data defined a homomorphism $\mathfrak{su}(\mathfrak{t}) \rightarrow Out(U(TYM))$.

An abstraction of the above observation is a homomorphism

$$h : \mathfrak{n} \rightarrow Out(C) \quad (36)$$

, where \mathfrak{n} is a (graded) Lie algebra, C is some (graded) associative algebra.

Under some assumptions, we can construct algebra B which enjoys the following list of properties:

1 There an isomorphism of linear spaces

$$\text{Sym}(\mathfrak{n}) \otimes C \stackrel{\mu}{\cong} B \quad (37)$$

,

2 C is a subalgebra of B , μ restricted on C is a homomorphism.

3 For $l, l_1, l_2 \in \mathfrak{n}$ $[\mu(l), \mu(C)] \subset \mu(C)$, $[\mu(l_1), \mu(l_2)] = \mu([l_1, l_2]) + \gamma(l_1, l_2) \in \mu(\mathfrak{n}) + C$. We require that it defines a homomorphism $\mathfrak{n} \rightarrow Out(C)$.

3' In case when C has an augmentation (as in case of universal enveloping) $\gamma(l_1, l_2) \in IC$, $[\mu(l), \mu(IC)] \subset \mu(IC)$

4 A map

$$l_1 \bullet \cdots \bullet l_n \otimes c \rightarrow \frac{1}{n!} \sum_{\sigma \in S_n} \pm \mu(l_{\sigma(1)}) \cdots \mu(l_{\sigma(n)}) \mu(c) \quad (38)$$

induces the isomorphism (37).

Definition 11 We say that the algebra belongs to the class $\mathbf{A}(\mathfrak{n}, C)$ ($\mathbf{A}'(\mathfrak{n}, C)$) if it satisfies assumptions **1,2,3,4** (**1,2,3',4**).

We provide a brief sketch of construction of $B \in \mathbf{A}(\mathfrak{n}, C)$ given a homomorphism h ((36)).

The map h can be used to construct an extension of Lie algebras

$$0 \rightarrow In(C) \rightarrow \tilde{\mathfrak{g}} \rightarrow \mathfrak{n} \rightarrow 0 \quad (39)$$

$In(C) = C/Z(C)$; $Z(C)$ is the center of C .

Denote by $L(C)$ a Lie algebra with linear space C and bracket defined by commutator in the associative algebra C .

In general it is not possible to lift (39) to an extension

$$0 \rightarrow L(C) \rightarrow \mathfrak{g} \rightarrow \mathfrak{n} \rightarrow 0 \quad (40)$$

with an obstruction in $H^3(\mathfrak{n}, Z(C))$.

Assumption 12 *Let us assume that $Z(C) = \mathbb{C}$ and there is an augmentation $C \rightarrow \mathbb{C}$. Then $L(C) = In(C) + Z(C)$ as a Lie algebra. This is satisfied if C is a free algebra.*

In such case we have a trivial lift of (39) to (40).

The universal enveloping $U(\mathfrak{g})$ contains $U(LC)$. There is a canonical homomorphism of associative algebras $U(LC) \rightarrow C$. Define $B = U(\mathfrak{g}) \underset{U(LC)}{\otimes} C$ - a free product of algebras. The isomorphism $B \cong \text{Sym}(\mathfrak{g}) \otimes C$ comes from Poincare-Birkhoff-Witt theorem.

Our experience shows that the scope of deformations of L in class $\mathbf{L}(\mathfrak{sush}(\alpha'), TYM(\alpha'))$ is a bit narrow. We expand it in the following

Definition 13 *We say that a deformation $U(L)_{\alpha'}$ of $U(L)$ is of $\mathbf{A}(\mathbf{A}')$ -type if it belongs to the class $\mathbf{A}(\mathfrak{sush}(\alpha'), U(TYM)(\alpha'))$ ($\mathbf{A}'(\mathfrak{sush}(\alpha'), U(TYM)(\alpha'))$).*

It is clear that any deformation of \mathbf{L} -type is of \mathbf{A} -type (take a universal enveloping algebra). The converse is not true.

In a study of deformations of type $\mathbf{A}(\mathfrak{sush}(\alpha'), U(TYM)(\alpha'))$ the theory of deformations of associative algebras might seem more relevant.

Suppose we have a deformation $L_{\alpha'} \in \mathbf{A}(\mathfrak{sush}(\alpha'), U(TYM)(\alpha'))$. Such deformations are governed by Hochschild cohomology (see section 7.1). There is a comparison result (28) which asserts that a deformation cocycle $\gamma(a, b)$ is completely determined by its values on $\Lambda^2(L) \subset U(L) \otimes U(L)$.

If $a, b \in TYM$, then $\gamma(a, b) \in U(TYM)$, because $U(TYM)(\alpha')$ is a subalgebra of $U(L)_{\alpha'}^1$.

¹As in Lie algebra case $U(TYM)$ is rigid

If $a \in L_1 + L_2, b \in TYM$ then $\gamma(a, b) \in U(TYM)$, because $[a, U(TYM)(\alpha')]_{\alpha'} \subset U(TYM)(\alpha')$.

If $a, b \in L_1 + L_2$ then $\gamma(a, b) \in U(TYM)$, because $L_1 + L_2$ generates homomorphism of $\mathfrak{su}\mathfrak{sh}$ into $Out(U(TYM)(\alpha'))$ and Lie algebra $\mathfrak{su}\mathfrak{sh}$ remains undeformed.

We summarize previous discussion in the following

Proposition 14 *Infinitesimal deformations of algebra $U(L)$ in a class of algebras $\mathbf{A}(\mathfrak{su}\mathfrak{sh}(\alpha'), U(TYM)(\alpha'))$ ($\mathbf{A}'(\mathfrak{su}\mathfrak{sh}(\alpha'), U(TYM)(\alpha'))$) are parametrized by points of a linear space $H^2(L, U(TYM))$ ($H^2(L, IU(TYM))$) of even degree. $H^2(L, U(TYM))^{Spin(10)}$ ($H^2(L, IU(TYM))^{Spin(10)}$) corresponds to $Spin(10)$ equivariant infinitesimal deformation.*

The algebra $U(L)_{\alpha'}$ contains a subalgebra $U(YM)_{\alpha'}$ generated by $\mu(\mathfrak{su}\mathfrak{sh}_2)(\alpha')$ -even part of $\mathfrak{su}\mathfrak{sh}$ and $U(TYM)(\alpha')$. The relations in $U(YM)_{\alpha'}$ have a form of (27), $r_j(\alpha'), s_\alpha(\alpha')$ have Taylor coefficients in associative algebra generated by (23). If we take a realization of this theory by connections and spinors, the perturbed part of YM-equation is a product of covariant derivatives of curvature and spinor fields.

Let us back up to the definition of deformations of type $\mathbf{L}(\mathfrak{su}\mathfrak{sh}(\alpha'), TYM(\alpha'))$. In this setup the deformed Lie algebra $L_{\alpha'}$ contains a subalgebra $YM_{\alpha'}$.

In our dictionary, where generators of YM correspond to the fields, cyclic words in generators correspond to Lagrangians. Indeed substituting fields for variables, taking trace multiplying on the volume form we get a Lagrangian. For this purpose is not necessary to use cyclic-ordinary words will do as good. However, representation of a Lagrangian as sum of ordinary words is redundant: trace has a cyclic symmetry as a result it picks up precisely cyclic words .

The algebraic analysis uses the following formalism. We have a short exact sequence

$$0 \rightarrow I \rightarrow \widetilde{YM} \xrightarrow{p} V \rightarrow 0 \quad (41)$$

The linear space V is equipped with a zero bracket. The map p projects v_1, \dots, v_{10} into a basis of V . The ideal I is generated by $F_{ij} = [v_i, v_j], \chi^\alpha$. As

a Lie algebra it is generated by symbols $[v_{(i_1}, [\dots [v_{i_k}, F_{ij}], [v_{(i_1}, [\dots [v_{i_k}, \chi^\alpha]$, where $()$ denote symmetrization, modulo relation $[v_i, F_{jk}] + [v_k, F_{ij}] + [v_j, F_{ki}] = 0$ (Bianchi identity). Denote mentioned space of generators by \widetilde{M} . The Lie algebra I is free. Making substitution (24) into an element n we recover $n(\nabla, \xi)$ from the introduction. The universal enveloping $U(I)$ has the same set of generators as I . The space of cyclic words ² in an alphabet defined by a basis of \widetilde{M} is $Cyc(U(I))$ is equal to $U(I)/[U(I), U(I)]$. The commutator $[U(I), U(I)]$ as linear space is generated by $ab - (-1)^{\deg(a)\deg(b)}ba$. The space $Cyc(U(I))$ has a homological interpretation. It is equal to $HH_0(U(I), U(I)) = H_0(I, U(I))$ (see sections 7.1 for definition HH_\bullet). The Lie algebra \widetilde{YM} acts on $Cyc(U(I))$ through abelian V . We have described above how to build a Lagrangian density $\mathcal{L}'(\nabla, \xi)$ out of element $\mathcal{L}' \in Cyc(U(I))$. It is clear that $(v_i \mathcal{L}')(\nabla, \xi) = \frac{\partial \mathcal{L}'(\nabla, \xi)}{\partial x_i}$. It means the action associated with $v_i \mathcal{L}'$ is trivial. To work directly with Lagrangians modulo full derivatives we shall replace $Cyc(U(I))$ by $Cyc(U(I))_V = H_0(\widetilde{YM}, U(I))$.

There is an additional subtlety when we deform a Lagrangian $\mathcal{L}(\nabla, \xi) \rightarrow \mathcal{L}(\nabla, \xi) + \alpha' \mathcal{L}'(\nabla, \xi)$ by $\mathcal{L}'(\nabla, \xi)$. If $\mathcal{L}'(\nabla, \xi) = 0$ on critical points of $\mathcal{L}(\nabla, \xi)$ then $\mathcal{L}(\nabla, \xi) + \alpha' \mathcal{L}'(\nabla, \xi)$ can be transformed to $\mathcal{L}(\nabla, \xi)$ by a field redefinition. In the algebraic language the space of on shell Lagrangians is equal to $H_0(\widetilde{YM}, U(I)/(\tilde{v}_m, \tilde{\chi}_\alpha)) = H_0(YM, U(TYM)) = H_0(YM, \text{Sym}(TYM))$, with $\tilde{v}_m, \tilde{\chi}_\alpha$ -defining relations of YM .

A simple way to deform YM is to deform original YM -LLagrangian. One of the conclusions in [14] was that Yang-Mills equations admit deformations of which are not Euler-Lagrange. It is natural then to impose the condition **Lg**:

Suppose we have a deformation $L_{\alpha'} \in \mathbf{L}(\mathfrak{su}(\eta)(\alpha'), TYM(\alpha'))$. Then the induced deformation of $YM_{\alpha'}$ should have the relations $\frac{\delta \mathcal{L}_{\alpha'}}{\delta v_i}, \frac{\delta \mathcal{L}_{\alpha'}}{\delta \chi^\alpha}$ coming from a Lagrangian $\mathcal{L}_{\alpha'}$.

Remark 15 *Abstract variational derivatives are uniquely defined by a property*

²Working in a graded setting we must take a grading into account while defining cyclic word

that

$$\begin{aligned} \left(\frac{\delta \mathcal{L}'}{\delta v_i} \right) (\nabla, \xi) &= \frac{\delta \mathcal{L}'(\nabla)}{\delta \nabla_i} \\ \left(\frac{\delta \mathcal{L}'}{\delta \chi^\alpha} \right) (\nabla, \xi) &= \frac{\delta \mathcal{L}'(\nabla)}{\delta \xi^\alpha} \end{aligned} \tag{42}$$

A similar definition can be given for the class of deformations $\mathbf{A}(\mathfrak{sush}(\alpha'), U(TYM)(\alpha'))$. Indeed the algebra $U(YM)_{\alpha'}$ in this setup is a subalgebra of $U(L)_{\alpha'}$ generated by $\mu(\mathfrak{sush}_2)$ and $U(TYM)_{\alpha'}$. The rest parallels to $\mathbf{L}(\mathfrak{sush}(\alpha'), TYM(\alpha'))$.

Let us say a few words about possible homological interpretation of a statement that \mathfrak{sush} closes on shell in N=1 D=10 Yang-Mills theory. Indeed the space of functions (functionals) on the space of fields as we know can be substituted by $Cyc(U(I))$ or $\text{Sym}(Cyc(U(I)))$ if we would like to work with multiple products of traces. The on shell functional are $H_0(TYM, U(TYM))$ ($\text{Sym}(H_0(TYM, U(TYM)))$). The Lie algebra L acts on TYM by commutators. It continues on $H_0(TYM, U(TYM))$ and factors through $\mathfrak{sush} = L/TYM$.

The same construction goes through for \mathbf{A} , \mathbf{L} deformations and infinitesimal deformations. The key moment is that $TYM, U(TYM)$ are rigid. Thus the condition of infinitesimal on shell closure of \mathfrak{sush} is automatically satisfied for infinitesimal deformations of \mathbf{A} , \mathbf{L} type.

3 Deformation complexes $U(TYM) \otimes S$ and $TYM \otimes S$.

In this section we start investigation of deformation cohomology $H^k(L, TYM)$ and $H^k(L, U(TYM))$, introduced in the section (2).

Section 7.1 offers some simplifications of deformation complexes. In the next few paragraphs we explain how results of section 7.1 can be adopted for our purposes.

Definition 16 Define a projective variety $\mathcal{Q} \subset \mathbf{P}^{15}$ by equations

$$r_i = \Gamma_{\alpha\beta}^i \lambda^\alpha \lambda^\beta = 0 \tag{43}$$

Proposition 17 *The algebra of homogeneous functions $S = \mathbb{C}[\lambda^1, \dots, \lambda^{16}]/(r_i)$ is Koszul (see [2]). By [15] and [16] $U(L) = S^!$*

Proposition 18 *The deformation cohomology $H^k(L, U(TYM))$ is equal to the k -th cohomology of the complexes $U(TYM) \otimes S$. The differential is a commutator with element*

$$e = \lambda^\alpha \theta_\alpha. \quad (44)$$

The cohomological grading coincides with the grading of S -factor. The total degree is preserved by d . The complex $U(TYM) \otimes S$ splits according to degree. A finer splitting can be achieved by identifying $U(TYM) = \bigoplus_{i \geq 0} \text{Sym}^i(TYM)$.

$$\text{Sym}^j(TYM)_n \otimes S_0 \rightarrow \text{Sym}^j(TYM)_{n+1} \otimes S_1 \rightarrow \dots \quad (45)$$

Proof. By remark (28) the deformation cohomology $H^k(L, U(TYM))$ is equal to $H^k(U(L), U(TYM))$. The bimodule $U(TYM)$ has left $U(L)$ -action induced by adjoint action of L . The right $U(L)$ -action is induced by the trivial L -action. A similar $U(L)$ -bimodule structure exists on $U(TYM)$. The propositions 32, (17) applied to $U(L)$ and our modules finish the proof. ■

Proposition 19 *Homology $H_k(L, U(TYM))$ are equal to the cohomology of the complex $U(TYM) \otimes S^*$. The space $S^* = \bigoplus_{n \geq 0} S_n^*$ is a bimodule dual to S . The differential is a commutator with element e ((44)). The homological degree coincides with the grading of S^* -factor. The complex $U(TYM) \otimes S^*$ also splits :*

$$\text{Sym}^j(TYM)_{3j} \otimes S_m^* \rightarrow \dots \rightarrow \text{Sym}^j(TYM)_{3j+m} \otimes S_0^* \quad (46)$$

Proof. Is similar to the proof of proposition (18). ■

The complex group $Spin(10, \mathbb{C})$ acts transitively on \mathcal{Q} ; the stable subgroup of a point is a parabolic subgroup P . To describe the Lie algebra \mathfrak{p} of P we notice that the Lie algebra $\mathfrak{so}(10, \mathbb{C})$ of $SO(10, \mathbb{C})$ can be identified with $\Lambda^2(V)$ (with the space of antisymmetric tensors ρ_{ab} where $a, b = 1, \dots, 10$). The vector

representation V of $SO(10, \mathbb{C})$ restricted to the group $GL(5, \mathbb{C}) \subset SO(10, \mathbb{C})$ is equivalent to the direct sum $W \oplus W^*$ of vector and covector representations of $GL(5, \mathbb{C})$. The Lie algebra of $SO(10, \mathbb{C})$ as vector space can be decomposed as $\Lambda^2(W) + \mathfrak{p}$ where $\mathfrak{p} = (W \otimes W^*) + \Lambda^2(W^*)$ is the Lie subalgebra of \mathfrak{p} . Using the language of generators we can say that the Lie algebra $\mathfrak{so}(10, \mathbb{C})$ is generated by skew-symmetric tensors m_{ab}, n^{ab} and by k_a^b where $a, b = 1, \dots, 5$. The subalgebra \mathfrak{p} is generated by k_a^b and n^{ab} . Corresponding commutation relations are

$$[m, m'] = [n, n'] = 0 \quad (47)$$

$$[m, n]_a^b = m_{ac} n^{cb} \quad (48)$$

$$[m, k]_{ab} = m_{ac} k_b^c + m_{cb} k_a^c \quad (49)$$

$$[n, k]_{ab} = n^{ac} k_c^b + n^{cb} k_c^a \quad (50)$$

The complex group $Spin(10)$ contains a two-sheet cover $\tilde{GL}(5)$ of $GL(5)$. Denote by W the fundamental representation of $GL(5)$. $\tilde{GL}(5)$ is the minimal cover on which $\det(g)^{\frac{1}{2}}, g \in \tilde{GL}_5$ becomes a single-valued representation. We denote it by $\det(W)^{\frac{1}{2}}$. We denote $\tilde{P} = \tilde{GL}(5) \ltimes \Lambda^2(W^*)$ the two-sheet cover of P . The space $W \otimes \det(W)^{-\frac{1}{2}}$ is \tilde{P} -representation.

Denote the bundle on \mathcal{Q} induced from $W \otimes \det(W)^{-\frac{1}{2}}$ by \mathcal{W} . The line bundle induced from $\det(W)^{\frac{1}{2}}$ by $\mathcal{O}(i)$. A notation $\mathcal{F} \otimes \mathcal{O}(i) = \mathcal{F}(i)$ is common in algebraic geometry.

The vector bundles $\Lambda^j \mathcal{W}(i)$ are $Spin(10)$ homogeneous. Thus the cohomology $H^k(\mathcal{Q}, \Lambda^j \mathcal{W}(i))$ are $Spin(10)$ -representations.

Proposition 20 *There is a long exact sequence connecting $H^k(L, U(TYM))$ and $H_k(L, U(TYM))$:*

$$\begin{aligned} \cdots \rightarrow H_{3-i, a-8}(L, \text{Sym}^j(TYM)) &\xrightarrow{\delta} H^{i, a}(L, \text{Sym}^j(TYM)) \rightarrow \\ \rightarrow H^{i+a-3j}(\mathcal{Q}, \Lambda^j(W)(3j-a)) &\xrightarrow{\iota} H_{2-i, a-8}(L, \text{Sym}^j(TYM)) \rightarrow \cdots \end{aligned} \quad (51)$$

Proof. See section (7.2). ■

Proposition 21 *Denote a sheaf of local holomorphic sections of $\Lambda^j \mathcal{W}(i)$ by the same symbol.*

$$i \geq 0$$

$$\begin{aligned}
H^0(\mathcal{Q}, \mathcal{W}(i+1)) &= [1, 0, 0, 0, i], & H^{10}(\mathcal{Q}, \mathcal{W}(-8-i)) &= [0, 0, 0, i, 1], \\
H^0(\mathcal{Q}, \Lambda^2 \mathcal{W}(2+i)) &= [0, 1, 0, 0, i], & H^{10}(\mathcal{Q}, \Lambda^2 \mathcal{W}(-8-i)) &= [0, 0, 1, i, 0], \\
H^9(\mathcal{Q}, \Lambda^2 \mathcal{W}(-6)) &= [0, 0, 0, 0, 0] \\
H^0(\mathcal{Q}, \Lambda^3 \mathcal{W}(3+i)) &= [0, 0, 1, 0, i], & H^{10}(\mathcal{Q}, \Lambda^3 \mathcal{W}(-7-i)) &= [0, 1, 0, i, 0], \\
H^1(\mathcal{Q}, \Lambda^3 \mathcal{W}(1)) &= [0, 0, 0, 0, 0] \\
H^0(\mathcal{Q}, \Lambda^4 \mathcal{W}(3+i)) &= [0, 0, 0, 1, i], & H^{10}(\mathcal{Q}, \Lambda^4 \mathcal{W}(-6-i)) &= [1, 0, 0, i, 0], \\
H^0(\mathcal{Q}, \Lambda^5 \mathcal{W}(3+i)) &= [0, 0, 0, 0, i], & H^{10}(\mathcal{Q}, \Lambda^5 \mathcal{W}(-5-i)) &= [0, 0, 0, i, 0],
\end{aligned} \tag{52}$$

denote $\mathcal{F}(t) = \sum_{i \geq 0} \dim(H^0(\mathcal{F}(i)))t^i$. Then

$$\begin{aligned}
\mathcal{O}(t) &= \frac{1 + 5t + 5t^2 + t^3}{(1-t)^{11}} \\
\mathcal{W}(t) &= \frac{10 + 34t + 16t^2}{(1-t)^{11}} \\
\Lambda^2 \mathcal{W}(t) &= \frac{45 + 65t + 11t^2 - t^3}{(1-t)^{11}} \\
\Lambda^3 \mathcal{W}(t) &= \frac{120 - 120t + 330t^2 - 462t^3 + 462t^4 - 330t^5 + 165t^6 - 55t^7 + 11t^8 - t^9}{(1-t)^{11}} \\
\Lambda^4 \mathcal{W}(t) &= \frac{16 + 34t + 10t^2}{(1-t)^{11}}
\end{aligned} \tag{53}$$

Proof. The proof reduces to a simple but tedious application of Borel-Weyl-Bott theorem, which was facilitated by a use of LiE program. ■

The Lie algebra TYM has a lowest graded component in the degree three (see (22)). It gives a nontrivial subspace in $H_0(L, TYM)$. Similarly the subspaces $\Lambda^j(L_3)$ isomorphically map into $H_{0,3j}(L, \text{Sym}^j(TYM))$ - subspace of degree $3j$.

Proposition 22 *The following maps and inclusions are isomorphisms for $i \geq 4$.*

$$\begin{aligned}
H^{10}(\mathcal{Q}, \Lambda^1(\mathcal{W})(-8-i)) &= [0, 0, 0, i, 1] \xrightarrow{\iota^1} H_{i,i+3}(L, TYM) \subset \\
&\subset H_i(L, TYM) \\
H^{10}(\mathcal{Q}, \Lambda^2(\mathcal{W})(-8-i)) &= [0, 0, 1, i, 0] \xrightarrow{\iota^2} H_{i,i+6}(L, \text{Sym}^2(TYM)) \subset \\
&\subset H_i(L, \text{Sym}^2(TYM)) \\
H^{10}(\mathcal{Q}, \Lambda^3(\mathcal{W})(-8-i)) &= [0, 1, 0, i+1, 0] \xrightarrow{\iota^3} H_{i,i+9}(L, \text{Sym}^3(TYM)) \subset \\
&\subset H_i(L, \text{Sym}^3(TYM)) \tag{54} \\
H^{10}(\mathcal{Q}, \Lambda^4(\mathcal{W})(-8-i)) &= [1, 0, 0, i+2, 0] \xrightarrow{\iota^4} H_{i,i+12}(L, \text{Sym}^4(TYM)) \subset \\
&\subset H_i(L, \text{Sym}^4(TYM)) \\
H^{10}(\mathcal{Q}, \Lambda^5(\mathcal{W})(-8-i)) &= [0, 0, 0, i+3, 0] \xrightarrow{\iota^5} H_{i,i+15}(L, \text{Sym}^5(TYM)) \subset \\
&\subset H_i(L, \text{Sym}^5(TYM)) \\
0 &= H_i(L, \text{Sym}^j(TYM)), \quad j \geq 6
\end{aligned}$$

Proof. Homology and cohomology of any Lie algebra are equal to zero in negative degree. $\Lambda^j(\mathcal{W}) = 0$ for $j \geq 6$. These observations together with long exact sequence (51) are sufficient for the proof. ■

4 Classification of deformation cocycles in

$$H^2(L, \text{Sym}(TYM))$$

We make a classification of deformation cocycles in $H^2(L, \text{Sym}(TYM))$ by its relation kernel and image of maps δ (51) and d_{dR}^L 89.

4.1 Exceptional classes that are not in the image of δ -type a .

These are the classes that span linear representations: $[0, 0, 0, 0, 2] \subset H^{2,-2}(L, \mathbb{C})$, $[1, 0, 0, 0, 1] \subset H^{2,1}(L, TYM)$, $[0, 1, 0, 0, 0] \subset H^{2,4}(L, \text{Sym}^2(TYM))$, which are

preimages of (95). There is also one sporadic (coming from higher cohomology): $c_{2,8} \in H^{2,8}(L, \text{Sym}^3(TYM))$ - a preimage of one of (96).

Proposition 23 *In the complex $S \otimes \text{Sym}^3(TYM)$ the class $c_{2,8}$ is represented by an element $\Gamma_{\alpha\beta}^{i_1, \dots, i_5} \Gamma_{\gamma\delta i_1 i_2 i_3} \lambda^\alpha \lambda^\beta \otimes \chi^\gamma \bullet \chi^\delta \bullet F_{i_4 i_5} \in S_2 \otimes \text{Sym}^3(TYM)_{10}$.*

Proof. It is easy to see that $\text{Sym}^3(TYM)_{10} = \Lambda^2(L_3) \otimes L_4$. The elements $c_{2,8}$ is the only $Spin(10)$ -invariant element in $S_2 \otimes \Lambda^2(L_3) \otimes L_4$. ■

Remark 24 *This deformation cocycle was analyzed in [1], [5] in connection with \mathcal{L}'_I .*

We leave as an exercise for the reader to check that all other classes belong to the image of δ . (Hint: use long exact sequence (51))

4.2 Classes that are in the image of δ .

These classes have their origin in homology group $H_1(L, \text{Sym}^j(TYM))$. In the section (7.4) we shall introduce an equivariant version of Connes differential

$$\dots \xrightarrow{d_{dR}^L} H_i(L, \text{Sym}^j(TYM)) \xrightarrow{d_{dR}^L} H_{i+1}(L, \text{Sym}^{j-1}(TYM)) \xrightarrow{d_{dR}^L} \dots \quad (55)$$

It satisfies $(d_{dR}^L)^2 = 0$. According to proposition (63) d_{dR}^L is zero on $H_1(L, \text{Sym}^j(TYM))$. It means that all elements of the later group are d_{dR}^L cocycles and can be classified with respect to d_{dR}^L .

Nontrivial d_{dR}^L cocycles in $H_1(L, \text{Sym}^j(TYM))$ -type b.

This linear space according to corollary (64) is a direct sum $A_j^1 + B_j^1$. The content of B_j^1 is written down in the lower part of (93). The elements of these representations are mapped to nontrivial deformation classes in $H^2(L, \text{Sym}(TYM))$ by the map δ .

The classes $\gamma_{1,4} \in H_1(L, TYM)$, $\gamma_{1,12} \in H_1(L, \text{Sym}^3(TYM))$ are the only $Spin(10)$ -invariants and are given by explicit formulas (94). Due to uniqueness of such classes and degrees counting we can identify the deformation $\delta\gamma_{1,4}$ with a deformation corresponding to infinitesimal Lagrangian \mathcal{L}'_{II} . The deformation $\delta\gamma_{1,12}$ corresponds to the Lagrangian \mathcal{L}'_{III} which yet to be constructed (some

information can be found in section (6)). The classes A_j^i according to (64) are mapped by δ to zero, therefore do not concern us.

d_{dR}^L **trivial cocycles in $H_1(L, \text{Sym}^j(TYM))$ - type c.**

These are elements in $H_1(L, \text{Sym}^j(TYM))$ of the form $d_{dR}^L \gamma$, where $\gamma \in H_0(L, \text{Sym}^{j+1}(TYM))$. We have an infinite-dimensional space of such elements.

Remark 25 *Any element of $\text{Sym}^{j+1}(TYM)$ produces (possibly zero) element in $H_0(L, \text{Sym}^{j+1}(TYM))$, for this we do not need to solve any equations, in contrast with elements of $H_1(L, \text{Sym}^{j+1}(TYM))$.*

We defer discussion of **Lg** properties of constructed deformation cocycles until the end of section (6). Briefly we may say that all such deformations are of **Lg** type.

5 A generating function for supersymmetric deformations

The groups which govern deformations $H^2(L, U(TYM))$ or $H^2(L, TYM)$ have a grading. This enables us to form generating functions:

$$\begin{aligned} a(t) &= \sum_{i \geq -2} \dim H^{2,i}(L, U(TYM)) t^i \\ l(t) &= \sum_{i \geq 1} \dim H^{2,i}(L, TYM) t^i \end{aligned} \tag{56}$$

In this section we procure formulas for $a(t)$ and $l(t)$. The section will lack proofs: the reader can use [14] as guide to recover the missed points.

The main ingredients of our formulas are:

$$\begin{aligned}
TYM(t) &= \sum_{i \geq 3} \dim TYM_i t^i = \sum_{i \geq 3} \frac{(-1)^i}{i} \times \\
&\times \sum_{kd=i} \mu(d) (11 + (-1)^{k-1} \{1 + \frac{1}{(2+\sqrt{3})^k} + \frac{1}{(2-\sqrt{3})^k}\}) t^i \\
\overline{HC}_0(U(TYM))(t) &= - \sum_{k \geq 1} \ln(1 - M((-1)^{k+1} t^k)) \frac{\psi(k)}{k} \\
U(TYM)(t) &= \sum_{i \geq 3} \dim U(TYM)_i t^i = \frac{1}{1 - M(t)} \\
1 - M(t) &= \frac{1 - 4t + t^2}{(1-t)^4(1-t^2)^5}
\end{aligned} \tag{57}$$

These have mostly a combinatorial origin.

$\mu(n)$ is the Mobius function. $\mu(n) = 0$ if n has repeated prime factors, $\mu(1) = 1$, $\mu(n) = (-1)^k$ if n is a product of k distinct primes.

$\psi(n)$ is the Euler function. is defined as the number of positive integers $\leq n$ that are relatively prime to n .

$$\begin{aligned}
p_0(t) &= \frac{1 - 5t + 5t^2 - t^3}{(1+t)^{11}} \\
p_1(t) &= t^3 \frac{16 - 34t + 10t^2}{(1+t)^{11}} - t^4 \\
p_2(t) &= t^6 \frac{120 + 120t + 330t^2 + 462t^3 + 462t^4 + 330t^5 + 165t^6 + 55t^7 + 11t^8 + t^9}{(1+t)^{11}} - 45t^8 \\
p_3(t) &= -t^8 \frac{45 - 65t + 11t^2 + t^3}{(1+t)^{11}} + 45t^8 + 10t^{10} - t^{12} - 144t^{11} \\
p_4(t) &= t^{10} \frac{10 - 34t + 16t^2}{(1+t)^{11}} - 10t^{10} + 144t^{11} + 16t^{13} - 126t^{14} \\
p_5(t) &= -t^{12} \frac{1 - 5t + 5t^2 - t^3}{(1+t)^{11}} + t^{12} - 16t^{13} + 126t^{14}
\end{aligned} \tag{58}$$

The last set of formulas is an appropriate modification of 53.

Finally

$$\begin{aligned}
a(t) &= \left(126t^{-2} + 144t + 45t^4 + t^8\right) + \left(t^{12} + 144t^{19} + t^{20} + 126t^{22}\right) + \\
&+ \left(-\frac{t^8(1-t)^{10}}{(1+t)^6} \overline{HC}_0(U(TYM))(t) - t^8 \left(\sum_{i=1}^5 ip_i(t)\right)\right) = a + b + c \quad (59) \\
l(t) &= \left(144t\right) + \left(211t^8\right) + \left(t^8 p_1(t) - t^8 p_0(t) TYM(t)\right) = a' + b' + c'
\end{aligned}$$

In $a(t)$ the summand a accounts for elements of type **a** from section (4.1), term b for type **b**, c for type **c**. The same interpretation holds for $l(t)$.

Finally we tabulated dimensions of $Spin(10)$ -invariants in $H^{2,deg}(L, \text{Sym}^p(TYM))$ in the following table. We use LiE program in our computations.

	4	8	12	16	20	24	28	32	<i>deg</i>
1			1		1	3	18	172	...
2							13	281	...
3		1			1	2	20	267	...
4							1	68	...
5							1	17	...
6									...
$\text{Sym}^p(TYM)$

(60)

This table is a superposition of three tables according to types introduced in section 4. The entries **a,b** tables were computed in 4. The table of type **c** is the biggest but the most regular. It was computed through Euler characteristic of $H_{\bullet,k}(L, \text{Sym}^i(TYM))$. The main ingredients in the computation are: the complex (55), propositions (22, 61), (62, 92). It is possible in principle to write a generating function for the numbers from the above table in terms of characters of few basic $Spin(10)$ representations and Adams operations.

6 Relation between ordinary and supersymmetric deformations

In this section we shall characterize a place of supersymmetric deformations among ordinary deformations. Ordinary (nonsupersymmetric) deformations

of YM are deformations of class $\mathbf{L}(V(\alpha'), TYM(\alpha'))$, $\mathbf{A}(V(\alpha'), U(TYM)(\alpha'))$ or a $Spin(10)$ -equivariant version, where V is an even part of $\mathfrak{su}\mathfrak{sh}$. As in the case of supersymmetric analogs such deformations are parametrized by $H^2(YM, TYM)$, $H^2(YM, U(TYM))$ or $Spin(10)$ -invariants of one of the spaces.

Since $U(TYM) = \text{Sym}(TYM)$ as YM -modules in our study we shall investigate relations between supersymmetric and ordinary cohomology with coefficients in $\text{Sym}^j(TYM)$.

There is an operation of restriction of chains

$$res : H^m(L, \text{Sym}^k(TYM)) \rightarrow H^m(YM, \text{Sym}^k(TYM)) \quad (61)$$

It is the map res what we shall concentrate on in this section. We shall characterize its kernel and image. In our study i is equal to 2.

We use a spectral sequence of extension $YM \subset L$. The E_1^{mn} -term is equal to

$$C^m(L/YM, H^n(YM, \text{Sym}^k(TYM))) \Rightarrow H^{m+n}(L, \text{Sym}^k(TYM)) \quad (62)$$

The classes which contribute to $H^2(L, \text{Sym}^k(TYM))$ are subquotients of

$$\begin{aligned} & C^0(L/YM, H^2(YM, \text{Sym}^k(TYM))) \\ & C^1(L/YM, H^1(YM, \text{Sym}^k(TYM))) \\ & C^2(L/YM, H^0(YM, \text{Sym}^k(TYM))) \end{aligned} \quad (63)$$

The classes of the last two linear spaces are mapped to zero under map res because of nontrivial L/YM -dependence. In proposition (84) we prove that $H^0(YM, \text{Sym}^k(TYM)) = 0$ for $k \geq 1$. Thus the last term makes a contribution to the kernel of res for $k = 0$, degree $i = -2$. The classes in $C^1(L/YM, H^1(YM, \text{Sym}^k(TYM)))$ by proposition (84) have degree $i = -3, -4$ for $k = 0$ and $i = 0, 1$ for $k = 1$. For $k \geq 2$ the corresponding linear space is zero. From long exact sequence (51) nontrivial classes in $H^{2,i}(L, \text{Sym}^k(TYM))$ for mentioned above degrees and values of k exist only for $k = 0, i = -2, k = 1, i = 1$. The group $H^{2,-2}(L, \mathbb{C})$ is isomorphic to $[0, 0, 0, 0, 2]$, $H^{2,1}(L, TYM)$ is isomorphic to $[1, 0, 0, 0, 1]$ as $Spin(10)$ -representation.

Proposition 26 *The kernel of map (61) ($m = 2$) is nontrivial only if $i = -2$ $k = 0$ and is equal to $[0, 0, 0, 0, 2]$; in degree $i = 1$ $k = 1$ and is equal to $[1, 0, 0, 0, 1]$*

The deformations which belong to $\text{Ker}(\text{res})$ of odd degree are unphysical. We can interpret $\text{Ker}(\text{res})$ as deformation of action of supersymmetries on YM , while the algebra YM is kept undeformed. It is easy to see that a deformation of degree -2 though changes L does not affect the adjoint action of $\theta_\alpha(\alpha')$ on $YM(\alpha')$. Thus this deformation can also be discarded.

Proposition 27 *The image of res in $H^2(YM, \text{Sym}^k(TYM))$ in $\text{deg} \geq 1$ can be characterized as classes invariant with respect to \mathfrak{susy} .*

Proof. We shall analyze higher differentials in this spectral sequence. There is only one possible nontrivial differential acting on

$$H^0(L/YM, H^2(YM, \text{Sym}^k(TYM)))$$

$$d_2 : H^0(L/YM, H^2(YM, \text{Sym}^k(TYM))) \rightarrow H^2(L/YM, H^1(YM, \text{Sym}^k(TYM))) \quad (64)$$

An interpretation of this map is the following. An element

$\gamma \in H^2(YM, \text{Sym}^k(TYM))$ defines infinitesimal deformation of YM . If it belongs to

$$H^0(L/YM, H^2(YM, \text{Sym}^k(TYM))) \subset H^2(YM, \text{Sym}^k(TYM))$$

then each supersymmetry can be modified to $\theta_\alpha + \alpha' \theta'_\alpha$ such that it defines a derivation of the infinitesimal deformation. A commutator of such two derivations defines an even derivation of YM , whose α' coefficient $\tau_{\alpha\beta}$ for fixed $\alpha\beta$ is an element of $H^1(YM, \text{Sym}^k(TYM))$ (recall a discussion in section (2) about derivations). The later group was computed in (64). It is nontrivial for $k = 0, 1$ and for $k = 1$ can be identified with \mathfrak{susy} . The whole object $\tau_{\alpha\beta}$ is an element of $H^2(L/YM, H^1(YM, \text{Sym}^k(TYM)))$. The condition that $\gamma \in \text{Ker} d_2$ is equivalent to a property that infinitesimally the action of an even translation receives a correction which can be eliminated by a field redefinition. It does not mean that we can do it simultaneously for all translations.

In the table below the reader may see degrees i of nonzero components of source and target (64) for different values of k .

k	$H^0(L/YM, H^2(YM, \text{Sym}^k(TYM)))$	$H^2(L/YM, H^1(YM, \text{Sym}^k(TYM)))$
0	$-6, -5$	$-5, -4$
1	$-2, -1, 0, \dots$	$-1, 0$
2	$1, 2, \dots$	
3	$4, 5, \dots$	
\dots	\dots	

Due to homogeneity with respect to \deg the map could be nontrivial only for $k = 0$ in degree -5 , for $k = 1$ in degree 0 . In both cases we do not have $Spin(10)$ -invariant deformations. ■

We need to remind to the reader (trivial from the point of view of the calculus of variations) relation between Lagrangians and deformations.

The first fact is that the Lie algebra YM is an algebra with Poincare duality (see [16]). It means that for any graded module N we have

$$H^{i,j}(YM, N) = H_{3-i,j+8}(YM, N), \quad (65)$$

the second index is the degree. It allows us to identify

$$H^2(YM, U(TYM)) \xrightarrow{P} H_1(YM, U(TYM)) = H_1(YM, \text{Sym}(TYM)) \quad (66)$$

The second fact is that there is a Connes differential defined in (7.4).

$$\dots \rightarrow H_i(YM, \text{Sym}^k(TYM)) \xrightarrow{d_{dR}^{YM}} H_{i+1}(YM, \text{Sym}^{k-1}(TYM)) \rightarrow \dots \quad (67)$$

It allows to map

$$H_{0,i}(YM, \text{Sym}^k(TYM)) \xrightarrow{var=P^{-1} \circ d_{dR}^{YM}} H^{2,i-8}(YM, \text{Sym}^{k-1}(TYM)) \quad (68)$$

It is easy to see that var has a simple variational interpretation. Yang-Mills theory has a Lagrangian $\mathcal{L}(\nabla, \xi)$. Pick an element $\mathcal{L}' \in H_0(YM, U(TYM))$ and define a deformed Lagrangian $\mathcal{L}(\nabla, \xi) + \alpha' \mathcal{L}'(\nabla, \xi)$. The deformed YM algebra according to remark (15) has relations $\frac{\delta \mathcal{L} + \alpha' \mathcal{L}'}{\delta v_i}, \frac{\delta \mathcal{L} + \alpha' \mathcal{L}'}{\delta \chi^\alpha}$. It is plausible and easy to check that cocycle of such deformation is $var(\mathcal{L}')$.

This enables us to translate into algebraic language condition **Lg**.

Suppose we have a deformation with a cocycle $c \in H^2(L, U(TYM))$. Then c is **Lg**-type if $res(c)$ belongs to the image of var and $var^{-1} \circ res(c)$ is the infinitesimal Lagrangian.

Remark (71) asserts that any cocycle in $H^2(L, U(TYM))^{Spin(10)}$ is of type **Lg** in the above sense.

Now we know everything to explain formulas (9) and (11). The factor t^8 corresponds to degree of Poincare duality map (65), because the later is involved in the definition of var . The terms which are subtracted from $a(t)$ define a generating function of $Kerres$. The relation of $l(t)$ and $\tilde{l}(t)$ is essentially the same.

We can explain relation of tables 13 and (60) along the same lines. The reader can see that the tables have the same entries, but a different numeration of rows and columns. The tables have the same entries because as we already know res is injective on $H^2(L, \text{Sym}(TYM))^{Spin(10)}$ and all $Spin(10)$ -invariants are in the image of var . The change of row numbering is due to increase $k \rightarrow k - 1$ in the definition of var in (68). The change in column numbering comes from two sources. Firstly, the map var changes degree deg by 8. Secondly in transition $deg \rightarrow deg_\alpha$ we use a formula (8).

In the introduction we made a claim that Lagrangian \mathcal{L}_{IV} can be constructed through the operator $\epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}}$. Let us elaborate on this claim. From our discussion at the end of section (2) about Lagrangians it should be clear that the right algebraization of $trn(\nabla, \xi)$ is an element of $Cyc(U(I))$. The last group maps to $H_0(L, U(TYM))$.

One of the ways to produce supersymmetry invariant infinitesimal Lagrangians would be to take an element $\mathcal{L}' \in H_0(L, \text{Sym}^k(TYM))$, map it to $H_1(L, \text{Sym}^{k-1}(TYM))$ via differential d_{dR}^L . Then use map δ to transport it to $H^2(L, \text{Sym}(TYM))$. Then by res to $H^2(YM, \text{Sym}(TYM))$. After that take a preimage var and end up in desired group $H_0(YM, \text{Sym}^k(TYM))$

One can do differently and use proposition (76). We know that $res \circ \delta \circ d_{dR}^L \mathcal{L}' = P^{-1} \circ d_{dR}^{YM} \circ P \circ res \circ \delta \mathcal{L}'$. This enables us to claim that $P \circ res \circ \delta \mathcal{L}'$ is an infinitesimal deformation of the Lagrangian corresponding to cocycle $\delta \circ d_{dR}^L \mathcal{L}'$.

Due to proposition (76) we may claim that $P \circ res \circ \delta \mathcal{L}' = \epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}} \mathcal{L}'$

7 Appendix

7.1 Cohomology of quadratic algebras

Material of this section is necessary for justification of reduction of big deformation complex $C(L, \text{Sym}(TYM))$ to a much smaller complex $S \otimes \text{Sym}(TYM)$. We use this in section (3)

It is well known that tangent space to the space of A_∞ deformations of an algebra A is governed by Hochschild cohomology $HH^\bullet(A, A)$. We define an object slightly more general - Hochschild cohomology with coefficients in a bimodule M . We denote it by $HH^\bullet(A, M)$. It is a cohomology of a complex

$$C^n(A, M) = \text{Hom}(A^{\otimes n}, M) \text{ with differential } d : C^n(A, M) \rightarrow C^{n+1}(A, M) \quad (69)$$

defined by the formula

$$\begin{aligned} d(c)(a_0, \dots, a_n) &= a_0 c(a_1, \dots, a_n) + \\ &+ \sum_{i=1}^n (-1)^i c(a_0, \dots, a_{i-1} a_i, \dots, a_n) + (-1)^{n+1} c(a_0, \dots, a_{n-1}) a_n \end{aligned} \quad (70)$$

If the algebra A and the module M are graded ($A = \bigoplus_{i=0}^\infty A_i$ with $A_0 = \mathbb{C}$, $M = \bigoplus_{i \in \mathbb{Z}} M_i$), then the complex $C^i(A, M)$ and cohomology are graded $C^{i,k} = \text{Hom}^k(A^{\otimes i}, M)$ by degree of a map. In the superscript $HH^{i,j}(A, M)$ the first is cohomological index, the second is the degree.

Remark 28 *This construction is related to the Lie algebra cohomology, described in section (2). Indeed if $A = U(\mathfrak{g})$, a bimodule M defines an adjoint \mathfrak{g} -module, which we denote by M^{ad} . The action of $l \in \mathfrak{g}$, $m \in M^{ad}$, then $lm = lm - ml$.*

According to [12] there is an isomorphism $HH^k(U(\mathfrak{g}), M) = H^k(\mathfrak{g}, M^{ad})$.

We also shall use a dual construction- homology groups $HH_i(A, M)$. Here it is brief description: Denote $C_n(A, M) = A^{\otimes n} \otimes M$ the groups which constitute

a complex $d : C_n(A, M) \rightarrow C_{n-1}(A, M)$ with a differential :

$$\begin{aligned} d(m \otimes a_1 \otimes \cdots \otimes a_n) &= ma_1 \otimes a_1 \otimes \cdots \otimes a_n + \\ &+ \sum_{i=1}^n (-1)^i m \otimes a_1 \otimes \cdots \otimes a_i a_{i+1} \otimes \cdots \otimes a_n \\ &+ (-1)^n a_n m \otimes a_1 \otimes \cdots \otimes a_{n-1} \end{aligned} \quad (71)$$

In our brief exposition of quadratic algebras we closely follow [17].

Definition 29 A graded algebra $A = \bigoplus_{i \geq 0} A_i$ is called a quadratic if $A_0 = \mathbb{C}$, $W = A_1$ generates A and all relations follow from quadratic relations $\sum_{i,j} r_{ij}^k \lambda^i \lambda^j = 0$ where $\lambda^1, \dots, \lambda^{\dim W}$ is a basis of W . The space of quadratic relations (the subspace of $W \otimes W$ spanned by $r_{ij}^1, r_{ij}^2, \dots$) will be denoted by R .

We see that $A_2 = W \otimes W / R$ and A is a quotient of free algebra (tensor algebra) $T(W)$ with respect to the ideal generated by R .

Definition 30 The dual quadratic algebra $A^!$ is defined as a quotient of a free algebra $T(W^*)$ generated by W^* (with a basis $\theta_1, \dots, \theta_{\dim W}$ dual to $\lambda^1, \dots, \lambda^{\dim W}$) by the ideal generated by $R^\perp \subset W^* \otimes W^*$ (here R^\perp stands for the subspace of $W^* \otimes W^* = (W \otimes W)^*$ that is orthogonal to $R \subset W \otimes W$).

There is an element e in the tensor product $A_1 \otimes A_1^! = W \otimes W^* = \text{End}(W)$. It corresponds to the identity element in $\text{End}(W)$. Equation $e^2 = 0$ is a direct corollary of orthogonality relations between R and R^\perp . The space $A^{!*}$ is dual bimodule. The left multiplication on e defines an operator

$$\dots \xrightarrow{d} A_i \otimes A_{n-i}^{!*} \xrightarrow{d} A_{i+1} \otimes A_{n-i-1}^{!*} \xrightarrow{d} \dots \quad (72)$$

It automatically satisfies $d^2 = 0$. The complex $A \otimes A^{!*}$ is called Koszul complex of A (the cohomological degree coincides with $A^{!*}$ grading)

Definition 31 A quadratic algebra A is called Koszul if $K_n(A)$ is acyclic for $n > 0$.

For Koszul algebra A there is a more economical complex than 69 suitable for computations of cohomology $HH^i(A, M)$.

Proposition 32 *Suppose A is a Koszul algebra and M is A -module. The cohomology of the complex $M \otimes A^!$ with a differential $d(a) = [e, a]$ is isomorphic to $HH^i(A, M)$. There is an appropriate refinement of this statement when M is graded.*

Proof.

A minor modification of $A \otimes A^{!*}$ is the complex $A \otimes A \otimes A^{!*}$. The differential is of a bicomplex $(A \otimes A \otimes A^{!*}, d_1, d_2)$. Let $e = \sum_{\alpha} \theta_{\alpha} \otimes \lambda^{\alpha}$. Denote $e_1 = \sum_{\alpha} \theta_{\alpha} \otimes 1 \otimes \lambda^{\alpha}$ and $e_2 = \sum_{\alpha} 1 \otimes \theta_{\alpha} \otimes \lambda^{\alpha}$, then d_1 is a left multiplication on e_1 and d_2 is a right multiplication on e_2 . Define a cohomological grading $A \otimes A \otimes A^{!*}$ by the degree of $A^{!*}$ factor.

Lemma 33 *Suppose A is a Koszul algebra. The complex $A \otimes A \otimes A^{!*}$ is acyclic in all degrees, but zero. Zero cohomology is equal to A .*

Proof. The proof is based on studying spectral sequence associated with a bicomplex $(A \otimes A \otimes A^{!*}, d_1, d_2)$, which collapses due to Proposition (31). ■

Recall that there is an interpretation of groups $HH^{\bullet}(A, M)$ through resolutions.

Definition 34 *Denote A^{op} an algebra with linear space of A and multiplication $a \times b = ba$.*

Any A -bimodule M becomes left $A \otimes A^{op}$ -module by the formula $(a \otimes b)m = amb$

Lemma 35 [4] *For any algebra A and a bimodule M there is an isomorphism of groups $HH^i(A, M) = Ext_{A \otimes A^{op}}^i(A, M)$. In the last group A is understood as $A \otimes A^{op}$ -module.*

This lemma may have the following interpretation: take any projective resolution of $A \otimes A^{op}$ -module A :

$$A \leftarrow P_0 \leftarrow P_1 \leftarrow \dots \leftarrow P_n \leftarrow \dots \quad (73)$$

Then the cohomology of the complex $Hom_{A \otimes A^{op}}(P_n, M)$ is equal to $HH^n(A, M)$ (the reader may consult [4] or [12] for the relevant definitions and details).

We can use a complex $A \otimes A \otimes A^{!*}$ as a resolution of A , where $P_n = A \otimes A \otimes A_n^{!*}$. It is a straightforward check that $\text{Hom}_{A \otimes A^{op}}(P_n, M) = A_n^{!*} \otimes M$. The differential $d(a)$ in this complex is a commutator of a with the canonical element e . If $a \otimes b \in M \otimes A^!$, then the homogeneity degree $l = \deg(a) - \deg(b)$ is preserved by the differential d . Then

$$M \otimes A^! = \bigoplus_l (M \otimes A^!)_l$$

-is direct sum of smaller complexes spanned by elements of homogeneity degree l which we denote by $(M \otimes A^!)_l$.

■

There is a version of this theory in homological setup.

Proposition 36 *Suppose A is a Koszul algebra and M is A -bimodule. Then homology groups $H_i(A, M)$ is isomorphic to cohomology of the complex $M \otimes A_n^{!*}$ with differential $d(a) = [e, a]$.*

Proof. Similar to the proof of proposition 32. Use flat $A \otimes A^{op}$ -resolution $A \otimes A \otimes A^{!*}$ of A . ■

7.2 Localization of $\text{Sym}^j(TYM) \otimes S$ and $\text{Sym}^j(TYM) \otimes S^*$

In this section we prove proposition (20).

Proposition 37 [16] *The cohomology $H^i(\mathcal{Q}, \mathcal{O}(n))$ is equal to S_n , if $i = 0$ and to S_k^* , if $n = -8 - k, i = 10, k \geq 0$. All other cohomology vanish*

Let $N = \bigoplus_{i \geq 0} N_i$ be a graded L -module. Consider a complex of vector bundles

$$N_0(l) \rightarrow N_1(l+1) \rightarrow \dots \rightarrow N_k(l+k) \rightarrow \dots \quad (74)$$

with differential defined by multiplication on element e from (44). In (74) N_n is understood as trivial vector bundle with a fiber N_n .

The $(\mathcal{O}(l)$ -twisted) localization of (45) or (46) is the complex (74) where $N_n = \text{Sym}^j(TYM)_n$. In such case by proposition (37) the complex of global sections of (74) for appropriate l and a shift is identical to (45), whereas the complex of tenth cohomology of (74) for suitable l and shift is equal to (46).

Definition 38 Denote the complex of vector bundles (74) by $\mathcal{N}(l)$. The bicomplex of Dolbeault of $(0, p)$ - forms with coefficients in $\mathcal{N}(l)$ we denote by $\Omega\mathcal{N}(l)$. It has two differentials: the original d and $\bar{\partial}$. The grading of the element in $\Omega^{0,p}N_i(l+i)$ is equal to $p+i$. The total complex we denote by $\Omega\mathcal{N}^\bullet(l)$.

There are two mappings- an embedding

$$i : H^0(\mathcal{Q}, \mathcal{N}(l))^\bullet \rightarrow \Omega\mathcal{N}^\bullet(l) \quad (75)$$

and projection

$$p : \Omega\mathcal{N}^\bullet(l) \rightarrow H^{10}(\mathcal{Q}, \mathcal{N}(l))^\bullet[-10] \quad (76)$$

Proposition 39 The map $p : \Omega\mathcal{N}^\bullet(l)/\text{Im}^\bullet(i) \rightarrow H^{10}(\mathcal{Q}, \mathcal{N}(l))^\bullet[-10]$ is a quasi-isomorphism.

Proof. Consider a spectral sequence of a bicomplex for $\Omega\mathcal{N}^\bullet(l)/\text{Im}^\bullet(i)$, in which we compute the cohomology of $\bar{\partial}$ first. The E_1 -term has nonzero entries equal to $E_1^{10,i} = H^{10}(\mathcal{Q}, \mathcal{O}(l+i)) \otimes N_i$, whence the proof. ■

Corollary 40 There is a long exact sequence of cohomology

$$\begin{aligned} \dots \rightarrow H^i(H^0(\mathcal{Q}, \mathcal{N}(l))) &\rightarrow H^i(\mathcal{Q}, \Omega\mathcal{N}(l)) \rightarrow \\ &\rightarrow H^{i-10}(H^{10}(\mathcal{Q}, \mathcal{N}(l))) \xrightarrow{\delta} H^{i+1}(H^0(\mathcal{Q}, \mathcal{N}(l))) \rightarrow \dots \end{aligned} \quad (77)$$

From now on we set N to be equal to $\text{Sym}^j(TYM)$. Our next goal is to compute the cohomology of a fiber of $\text{Sym}^j(\mathcal{T}\mathcal{Y}\mathcal{M})^\bullet(l)$ over a point $x \in \mathcal{Q}$. Due to $Spin(10)$ homogeneity, fibers over different points are isomorphic. This will enable us to compute $H^i(\mathcal{Q}, \Omega\text{Sym}^j(\mathcal{T}\mathcal{Y}\mathcal{M})(l))$ completely. To do this we need to remind some facts about the manifold \mathcal{Q} .

It is good deal is known about \mathcal{Q} (see [15] for list of properties and references.)

The main properties of \mathcal{Q} are :

1. The complex dimension of \mathcal{Q} is equal to 10.
2. As a real homogeneous manifold \mathcal{Q} is equal to $SO(10, \mathbb{R})/U(5)$.

It follows from this that \mathcal{Q} is smooth. The affine cone, defined by equation (43), is also smooth away from the apex.

If we choose λ_0^α - a solution to (43), then $\theta = \sum_\alpha \lambda_0^\alpha \theta_\alpha$ satisfies $[\theta, \theta] = 2\theta^2 = 0$. It can be used to define a differential d on L and on $U(L)$ by the formula $d(a) = [\theta, a]$. A vector λ_0^α is coordinates of point x on the cone $C\mathcal{Q}$.

One can define a one-dimensional S -bimodule \mathbb{C}_x by specialization at $x \in C\mathcal{Q}$ with coordinates λ_0^α . To emphasize x -dependence of the differential d we denote it by d_x .

Proposition 41 $HH^i(S, \mathbb{C}_x) = H^i(U(L), d_x)$

Proof. This is a direct application of proposition 32, where $M = \mathbb{C}_x$. ■

Proposition 42 *Suppose A is a ring of algebraic functions on affine algebraic variety. \mathbb{C}_x is a one-dimensional bimodule, corresponding to a smooth point x . Then $HH^i(A, \mathbb{C}_x) = \Lambda^i(T_x)$, where T_x is the tangent space at x .*

Proof. This is a weak form of Hochschild-Kostant-Rosenberg theorem. ■

Corollary 43 $HH^i(S, \mathbb{C}_x) = H^i(U(L), d_x) = \Lambda^i(T_x)$, where T_x is the tangent space to $C\mathcal{Q}$ at the point x , $x \neq 0$.

Remark 44 *This can be interpreted in elementary terms. Suppose we would like to deform d_x by deforming x . The condition*

$$\Gamma_{\alpha\beta}^i \lambda_0^\alpha \xi^\beta = 0 \quad i = 1, \dots, 10 \quad (78)$$

is the condition that a vector with coordinates ξ^α is tangent to $C\mathcal{Q}$ at x . This is precisely the condition that $[e_x, g] = 0$, where $g = \sum_\alpha \xi^\alpha \theta_\alpha$. Multiplicative structure in $U(L)$ allows to multiply different g 's, spanning the exterior algebra $\Lambda^i(T_x)$. A less trivial fact established in (43) is that this is how we can generate all the cohomology.

Corollary 45 $H^i(\text{Sym}^i(L), d_x) = \Lambda^i(T_x)$ and all other cohomology equal to zero.

Proof. By Poincare-Birkhoff-Witt theorem $(U(L), d_x) = (\text{Sym}(L), d_x)$. It means that we have a direct sum decomposition of subcomplexes $(\text{Sym}(L), d_x) = \bigoplus_{n \geq 0} (\text{Sym}^n(L), d_x)$. In the previous paragraph we established that $T_x \subset H^1(L, d_x)$. By construction $\Lambda^i(T_x) \subset H^i(\text{Sym}^i(L), d_x)$. By (43) $\Lambda(T_x)$ exhaust all the cohomology. ■

Our next goal is the computation of $H(\text{Sym}^j(TYM), d_x)$.

For this is useful to know $Spin(10)$ -representation content of the Lie algebra L .

A general algorithm to determine isotopic component in L is the following:

1. We know the representation content of homogeneous components of the algebra $S = \bigoplus_{n \geq 0} S_n$. The spinor representation S_1 has the highest weight $[0, 0, 0, 0, 1]$. The n -th component S_n (by Borel-Weyl-Bott theorem) is an irreducible representation of weight $[0, 0, 0, 0, n]$. Koszul property of S tells us that there is a series of acyclic complexes (differential is a left multiplication on e .)

$$U(L)_n \otimes S_0^* \leftarrow U(L)_{n-1} \otimes S_1^* \leftarrow \cdots \leftarrow U(L)_0 \otimes S_n^* \quad (79)$$

They can be used to compute $Spin(10)$ -representation in $U(L)_n$ inductively.

We can assume now that we know the content of all $U(L)_n$. The main formula which enables to recover the content of L_n from $U(L)_n$ is the Poincare-Birkhoff-Witt isomorphism $U(L) = \text{Sym}(L)$. It can be used to compute L_n by induction. Such computations are greatly facilitated by the use of LiE program. The low graded components of L are

$$\begin{aligned} L_1 &= [0, 0, 0, 1, 0] \\ L_2 &= [1, 0, 0, 0, 0] \\ L_3 &= [0, 0, 0, 0, 1] \\ L_4 &= [0, 1, 0, 0, 0] \\ &\dots \end{aligned} \quad (80)$$

Proposition 46 *There are the following identifications of \widetilde{GL}_5 -representations :*

$$\begin{aligned}
L_1 &= (\mathbb{C} + \Lambda^2(W) + \Lambda^4(W)) \otimes \det(W)^{-\frac{1}{2}} \\
L_2 &= W + W^* \\
L_3 &= (\mathbb{C} + \Lambda^2(W^*) + \Lambda^4(W^*)) \otimes \det(W)^{\frac{1}{2}} = \\
&= (\mathbb{C} + W \otimes \det(W)^{-1} + \Lambda^3(W)^{-1}) \otimes \det(W)^{\frac{1}{2}} \\
L_4 &= \Lambda^2(W) + \Lambda^2(W^*) + W \otimes W^* \\
&\dots
\end{aligned} \tag{81}$$

Proof. Use a realization of spinor representation in even (odd) exterior powers of W (Fock representation) described in [3]. ■

Proposition 47 $H^3(TYM, d_x) = W \otimes \det(W)^{-\frac{1}{2}}$, all other cohomology vanish

Proof.

It not hard to describe the action of the differential d_x on L_n for small n , where the point x is invariant with respect to SL_5 . We describe the differential d_x explicitly using decomposition (81). We included only those isotopic components on which d_x is not zero. On such components the differential is uniquely defined up to a scalar factor

$$\begin{aligned}
L_1 &\supset \Lambda^2(W) \otimes \det(W)^{\frac{1}{2}} = W^* \otimes \det(W)^{\frac{1}{2}} \xrightarrow{d_x} W^* \subset L_2 \\
L_2 &\supset W \xrightarrow{d_x} W \otimes \det(W)^{-\frac{1}{2}} \subset L_3 \\
L_3 &\supset \det(W)^{\frac{1}{2}} + \Lambda^3(W) \otimes \det(W)^{-\frac{1}{2}} \rightarrow \mathbb{C} + \Lambda^2(W^*) \subset W \otimes W^* + \Lambda^2(W^*)
\end{aligned} \tag{82}$$

The complex $L_1/T_x \rightarrow L_2 \rightarrow \dots$ is acyclic. If we truncate L_1/T_x and L_2 terms, the resulting complex will have cohomology equal to $d_x(L_2) = W \otimes \det(W)^{-\frac{1}{2}}$. ■

Corollary 48 *The complex $(\text{Sym}^j(TYM), d_x)$ has cohomology in degree $3i$ equal to $\Lambda^i(W \otimes \det(W)^{-\frac{1}{2}})$.*

Proof. Similar to the proof of corollary (45). ■

To find cohomology of $(\mathrm{Sym}^j(TYM) \otimes S, d)$ we plan study linear spaces $\Lambda^j(W_x \otimes \det(W_x)^{-\frac{1}{2}})$ as a family $x \in C\mathcal{Q} \setminus \{0\}$.

There is a natural identification of linear spaces $\Lambda^j(W_x \otimes \det(W_x)^{-\frac{1}{2}})$ and $\Lambda^j(W_{\lambda x} \otimes \det(W_{\lambda x})^{-\frac{1}{2}})$ $\lambda \in \mathbb{C}^\times$. From this we conclude that the family $W_x \otimes \det(W_x)^{-\frac{1}{2}}$ can be pushed to a vector bundle \mathcal{W} on \mathcal{Q} .

Proposition 49 $H^i(\Omega \mathrm{Sym}^j \mathcal{T}\mathcal{Y}\mathcal{M}(l)) = H^{i-3j}(\mathcal{Q}, \Lambda^j \mathcal{W}(l))$

Proof. In corollary (45) we found that local (at a point) cohomology of $\mathrm{Sym}^j \mathcal{T}\mathcal{Y}\mathcal{M}(l)$ is equal to $\Lambda^j(\mathcal{W})(l)[-3j]$. In fact the later vector bundle is a subbundle of $\mathrm{Sym}^j \mathcal{T}\mathcal{Y}\mathcal{M}(l)$. Thus the embedding $\Omega \Lambda^j(\mathcal{W})(l) \rightarrow \Omega \mathrm{Sym}^j \mathcal{T}\mathcal{Y}\mathcal{M}(l)$ is a quasiisomorphism. ■

Proof. of Proposition (51) One simply has to make necessary shift in cohomological degree and grading when identifies $H_{i,k}(L, \mathrm{Sym}^j(TYM))$ and $H^\bullet(H^{10}(\mathcal{Q}, \mathrm{Sym}^j(TYM)(l)))$. The same applies to cohomology. Then use proposition (49) and (40). ■

Proof. of lemma (58) A plan is to use multiplicative action of $H^i(\mathfrak{sush}, \mathbb{C})$ on $H_n(L, \mathrm{Sym}^j(TYM))$. The action of $H^i(\mathfrak{sush}, \mathbb{C})$ factors through the action of $H^i(L, \mathbb{C})$

The maps ι_j in (54) induced by tautological inclusion of sheaves $\Lambda^j \mathcal{W}(-8-i) \rightarrow \Lambda^j(L_3)(-8-i) \subset \mathrm{Sym}^j(TYM)(-8-i)$. We claim that ι_j are embedding. Let $p : \mathcal{F} \rightarrow \mathcal{Q}$ be a canonical projection from the full flags.

The idea is to decompose $\Lambda^j(L_3)$ into $Spin(10)$ -irreducible components. Pick one of them, denoted by A , such that $\Lambda^j \mathcal{W}(-8-i) \subset A(-8-i)$. Using multiplication on sections $H^0(\mathcal{Q}, \mathcal{O}(i))$ one can reduce the check to the case $i = 0$. The Serre dual statement is a claim that $A^* \rightarrow \Lambda^j \mathcal{W}^*$ induces isomorphism on global sections. Such diagram of sheaves is a pushforward of a diagram $A^* \rightarrow \mathcal{O}(D)^3$ on the full flags, where $p_* \mathcal{O}(D_j) = \Lambda^j \mathcal{W}^*$, $p_* A^* = A^*$. On the full flags it becomes the statement of classical Borel-Weyl theory that there is an isomorphism $A^* = H^0(\mathcal{F}, \mathcal{O}(D_j))$. In particular the map is an embedding for $i = 0$. The linear subspace $A \subset \Lambda^j(L_3)$ embeds into homology.

³ D is a divisor

We can reinterpret multiplications on sections $H^0(\mathcal{Q}, \mathcal{O}(i))$ as multiplication on elements of $H^i(\mathfrak{su}\mathfrak{s}\mathfrak{h}, \mathbb{C})$.

■

7.3 Homology of algebra $\mathfrak{su}\mathfrak{s}\mathfrak{h}$.

In this section we compute cohomology of $\mathfrak{su}\mathfrak{s}\mathfrak{h}$ in dimension 10 with trivial coefficients.

The complex that computes Lie algebra homology is equal to a space of (super)-polynomial functions $C_\bullet(\mathfrak{su}\mathfrak{s}\mathfrak{h}, \mathbb{C}) = \{f(u^1, \dots, u^{16}, \tau^1, \dots, \tau^{10})\}$ where u^α has degree $Kdeg$ zero, τ^i - one. The differential is given by the formula $\Gamma_i^{\alpha\beta} \xi^i \frac{\partial^2}{\partial u^\alpha \partial u^\beta}$. It is convenient to introduce a finer bigrading on $C_\bullet(\mathfrak{su}\mathfrak{s}\mathfrak{h}, \mathbb{C})$ by setting $deg(u^\alpha) = 1$ and $deg\xi^i = 2$. The complex $C_\bullet(\mathfrak{su}\mathfrak{s}\mathfrak{h}, \mathbb{C})$ then decomposes into a direct sum of complexes of elements of different degree deg .

deg	\dots	\dots	\dots	\dots	\dots	\dots
4	$S^4(S)$	\rightarrow	$S^2(S) \otimes \Lambda^1(V)$	\rightarrow	$\Lambda^2(V)$	\dots
3	$S^3(S)$	\rightarrow	$S^1(S) \otimes \Lambda^1(V)$			
2	$S^2(S)$	\rightarrow	$\Lambda^1(V)$			
1	$S^1(S)$					
0	\mathbb{C}					
	0		1		2	$Kdeg$

(83)

Proposition 50 *In the table below you can find the representation content of*

the homology groups $H_i(\mathfrak{su}\mathfrak{h}, \mathbb{C})$

<i>deg</i>	
15	[0, 0, 0, 15, 0]	[0, 0, 0, 12, 1]	[0, 0, 1, 9, 0]	[0, 1, 0, 7, 0]	[1, 0, 0, 5, 0]	[0, 0, 0, 3, 0]	
14	[0, 0, 0, 14, 0]	[0, 0, 0, 11, 1]	[0, 0, 1, 8, 0]	[0, 1, 0, 6, 0]	[1, 0, 0, 4, 0]	[0, 0, 0, 2, 0]	
13	[0, 0, 0, 13, 0]	[0, 0, 0, 10, 1]	[0, 0, 1, 7, 0]	[0, 1, 0, 5, 0]	[1, 0, 0, 3, 0]	[0, 0, 0, 1, 0]	
12	[0, 0, 0, 12, 0]	[0, 0, 0, 9, 1]	[0, 0, 1, 6, 0]	[0, 1, 0, 4, 0]	[1, 0, 0, 2, 0]	[0, 0, 0, 0, 0]	
11	[0, 0, 0, 11, 0]	[0, 0, 0, 8, 1]	[0, 0, 1, 5, 0]	[0, 1, 0, 3, 0]	[1, 0, 0, 1, 0]		
10	[0, 0, 0, 10, 0]	[0, 0, 0, 7, 1]	[0, 0, 1, 4, 0]	[0, 1, 0, 2, 0]	[1, 0, 0, 0, 0]		
9	[0, 0, 0, 9, 0]	[0, 0, 0, 6, 1]	[0, 0, 1, 3, 0]	[0, 1, 0, 1, 0]			
8	[0, 0, 0, 8, 0]	[0, 0, 0, 5, 1]	[0, 0, 1, 2, 0]	[0, 1, 0, 0, 0]			
7	[0, 0, 0, 7, 0]	[0, 0, 0, 4, 1]	[0, 0, 1, 1, 0]				
6	[0, 0, 0, 6, 0]	[0, 0, 0, 3, 1]	[0, 0, 1, 0, 0]				
5	[0, 0, 0, 5, 0]	[0, 0, 0, 2, 1]					
4	[0, 0, 0, 4, 0]	[0, 0, 0, 1, 1] + [0, 0, 0, 0, 0]					
3	[0, 0, 0, 3, 0]	[0, 0, 0, 0, 1]					
2	[0, 0, 0, 2, 0]						
1	[0, 0, 0, 1, 0]						
0	[0, 0, 0, 0, 0]						
	0	1	2	3	4	5	<i>Kdeg</i>

(84)

To find the homological dimension \dim of a class $\alpha \in H_i(\mathfrak{su}\mathfrak{h}, \mathbb{C})$ of the bidegree $\deg, Kdeg$, one needs to use a formula $i = \deg - Kdeg$. Reviewing the table we see that there are nontrivial $\mathfrak{so}(10)$ invariant classes in degrees 3 and 7.

Remark 51 The representation content of cohomology group can be read off from the table (84) by replacing a cell entry $[w_1, w_2, w_3, w_4, w_5]$ by $[w_1, w_2, w_3, w_5, w_4]$.

Proof. Rather than giving a full proof of the statements which is based on fairly standard technique we hint the main points.

Instead of homology, we compute cohomology. We complete the algebra $\mathbb{C}[u^1, \dots, u^{16}] \otimes \Lambda[\tau_1, \dots, \tau_{10}]$, by the ideal generated by $\Gamma_{\alpha\beta}^i u^\alpha u^\beta$. Due to the grading, preserved by the differential this operation does a completion of cohomology. The completed ring can be interpreted as a ring of homogeneous functions in a formal neighborhood of $\mathcal{Q} \times \mathbb{C}^{0|10} \subset \mathbf{P}^{15} \times \mathbb{C}^{0|10}$. Then we localize the complex on this neighborhood. The computation is based on spectral

sequences of hypercohomology of the localized complex. ■

7.4 Equivariant ordinary and cyclic homology.

For any Lie algebra \mathfrak{g} and a module N , the cochains $C^\bullet(\mathfrak{g}, N)$ is a module over differential graded algebra $C^\bullet(\mathfrak{g}, \mathbb{C})$. Similarly $C_\bullet(\mathfrak{g}, \mathbb{C})$ is a coalgebra and $C_\bullet(\mathfrak{g}, N)$ is a comodule over it. The groups $C_n(\mathfrak{g}, \mathbb{C})$ and $C^n(\mathfrak{g}, \mathbb{C})$ are dual with a pairing $\langle a, b \rangle$. Denote $\Delta(n) = \sum_i b_i \otimes n_i$ -the diagonal of $n \in C_n(\mathfrak{g}, N)$. This enables us to define an action $C^i(\mathfrak{g}, \mathbb{C}) \otimes C_n(\mathfrak{g}, N) \rightarrow C_{n-i}(\mathfrak{g}, N)$ by the formula $an = \sum_i \langle a, b_i \rangle n_i$.

It is easy to see that the action is compatible with differentials and induces action of $H^i(\mathfrak{g}, \mathbb{C})$ on $H_n(\mathfrak{g}, N)$.

There is one more related homology theory-cyclic homology. According to [10] in case of universal enveloping algebras there is a different way (different than a standard which is due to Connes [7]) to define cyclic homology which we adopt in this paper.

Definition 52 *Suppose a complex C is equipped with two differential b, B , $\deg B = 1$, $\deg b = -1$. The differentials satisfy $b^2 = B^2 = bB + Bb = 0$. Then C is called a mixed complex.*

For a mixed complex C define a bicomplex

$$\begin{array}{ccccc}
 \cdots & & \cdots & & \cdots \\
 \downarrow b & & \downarrow b & & \downarrow b \\
 C^2 & \xleftarrow{B} & C^1 & \xleftarrow{B} & C^0 \\
 \downarrow b & & \downarrow b & & \\
 C^1 & \xleftarrow{B} & C^0 & & \\
 \downarrow b & & & & \\
 C^0 & & & &
 \end{array}$$

Denote the total complex of the above bicomplex by $ToT(C, B, b)$. One can extend this to infinity to the left

$$\begin{array}{ccccccc}
& & \dots & & \dots & & \dots \\
& & \downarrow b & & \downarrow b & & \downarrow b \\
\dots & \xleftarrow{B} & C^2 & \xleftarrow{B} & C^1 & \xleftarrow{B} & C^0 \\
& & \downarrow b & & \downarrow b & & \\
\dots & \xleftarrow{B} & C^1 & \xleftarrow{B} & C^0 & & \\
& & \downarrow b & & & & \\
\dots & \xleftarrow{B} & C^0 & & & & \\
& & \dots & & & & \\
\hline
& & 0 & & 1 & & 2
\end{array}$$

Denote the resulting total complex by $ToT^{per}(C, B, b)$. We allow infinite sums in $ToT^{per}(C, B, b)$.

Similarly define a bicomplex with zero entries to the right from the zero column. Denote the complex by $ToT^-(C, B, b)$

Fix a Lie algebra \mathfrak{g} . On polynomial forms $\Omega(\mathfrak{g}^*)$ of coadjoint representations of \mathfrak{g} there are two differentials. The first one is the de Rham differential d_{dR} . The linear space of coadjoint representation is a Poisson manifold via Kirillov bracket $\{a, b\}$. The second differential is defined by the formula

$$\begin{aligned}
d(a_0 da_1 \dots da_n) &= \sum_{i=1}^n (-1)^i \{a_i, a_0\} da_1 \dots \widehat{da_i} \dots da_n + \\
&+ \sum_{i=1}^n (-1)^{i+j-1} a_0 d\{a_i, a_j\} da_1 \dots \widehat{da_i} \dots \widehat{da_j} \dots da_n
\end{aligned} \tag{85}$$

Suppose we have an extension of Lie algebras

$$0 \rightarrow \mathfrak{l} \rightarrow \mathfrak{g} \rightarrow \mathfrak{n} \rightarrow 0 \tag{86}$$

Define $\Omega^{n^*}(\mathfrak{g}^*)$ as polynomial differential forms on \mathfrak{g}^* , which are invariant with respect to \mathfrak{n}^* -translations. It is easy to see that $\Omega^{n^*}(\mathfrak{g}^*)$ is closed under differentials d_{dR} and d

Proposition 53 $\Omega(\mathfrak{g}^*)$ and $\Omega^{n^*}(\mathfrak{g}^*)$ are mixed complexes, with $B = d_{dR}$, $b = d$.

Definition 54 $HC_n(U(\mathfrak{g})) = H_n(ToT(\Omega(\mathfrak{g}^*), d_{dR}, d)),$

$$HC_n(\mathfrak{g}, U(\mathfrak{l})) = H_n(ToT(\Omega^{n^*}(\mathfrak{g}^*), d_{dR}, d)).$$

$$HC_n^{per}(U(\mathfrak{g})) = H_n(ToT^{per}(\Omega(\mathfrak{g}^*), d_{dR}, d)),$$

$$HC_n^{per}(\mathfrak{g}, U(\mathfrak{l})) = H_n(ToT^{per}(\Omega^{n^*}(\mathfrak{g}^*), d_{dR}, d))$$

$$\begin{aligned}
HC_n^-(U(\mathfrak{g})) &= H_n(ToT^-(\Omega(\mathfrak{g}^*), d_{dR}, d)), \\
HC_n^-(\mathfrak{g}, U(\mathfrak{l})) &= H_n(ToT^-(\Omega^{\mathfrak{n}^*}(\mathfrak{g}^*), d_{dR}, d))
\end{aligned}$$

The following long exact sequence easily follows from definitions

$$\begin{aligned}
\cdots \rightarrow HC_{n-1}(\mathfrak{g}, U(\mathfrak{l})) &\rightarrow HC_n^-(\mathfrak{g}, U(\mathfrak{l})) \rightarrow HC_n^{per}(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \\
&\rightarrow HC_{n-2}(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \cdots \\
\cdots \rightarrow HC_{n+2}^-(\mathfrak{g}, U(\mathfrak{l})) &\rightarrow HC_n^-(\mathfrak{g}, U(\mathfrak{l})) \rightarrow HC_n(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \\
&\rightarrow HC_{n+1}(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \cdots
\end{aligned} \tag{87}$$

$$\begin{aligned}
\cdots \rightarrow HC_{n-1}(\mathfrak{g}, U(\mathfrak{l})) &\rightarrow H_n(\mathfrak{g}, U(\mathfrak{l})) \rightarrow HC_n(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \\
&\rightarrow HC_{n-2}(\mathfrak{g}, U(\mathfrak{l})) \rightarrow \cdots
\end{aligned}$$

Consider a complex :

$$\begin{array}{ccccc}
\cdots & & \cdots & & \cdots \\
\downarrow d & & \downarrow d & & \downarrow d \\
\text{Sym}^0(\mathfrak{l}) \otimes \Lambda^2(\mathfrak{g}) & \xrightarrow{d_{dR}} & \text{Sym}^1(\mathfrak{l}) \otimes \Lambda^1(\mathfrak{g}) & \xrightarrow{d_{dR}} & \text{Sym}^2(\mathfrak{l}) \otimes \Lambda^0(\mathfrak{g}) \\
\downarrow d & & \downarrow d & & \\
\text{Sym}^0(\mathfrak{l}) \otimes \Lambda^1(\mathfrak{g}) & \xrightarrow{d_{dR}} & \text{Sym}^1(\mathfrak{l}) \otimes \Lambda^0(\mathfrak{g}) & & \\
\downarrow d & & & & \\
\text{Sym}^0(\mathfrak{l}) \otimes \Lambda^0(\mathfrak{g}) & & & &
\end{array} \tag{88}$$

Proposition 55 *The cohomology of the total complex 88 is equal to $H_\bullet(\mathfrak{n}, \mathbb{C})$.*

Proof. To prove we use the spectral sequence of the bicomplex which E_1 term is equal to the cohomology of 88 with respect to d_{dR} . Choosing some linear splitting of (86), we identify $\Lambda^n(\mathfrak{g}) = \bigoplus_{i+j=n} \Lambda^i(\mathfrak{n}) \otimes \Lambda^j(\mathfrak{l})$. The horizontal rows are equal to direct sum of homogeneous components of the de Rham complex on \mathfrak{l}^* with coefficients in $\Lambda^i(\mathfrak{n})$ for various i . Due to acyclicity of the de Rham complex of \mathfrak{l}^* the cohomology of n -th row is equal to $\Lambda^n(\mathfrak{n})$. These are located in the first column. It is obvious that the vertical differential becomes the standard homology differential in $C_n(\mathfrak{n}, \mathbb{C}) = \Lambda^n(\mathfrak{n})$, described in (32).

The spectral sequence collapses in E_2 -term due to dimension reasons. ■

Corollary 56 $HC_n^{per}(\mathfrak{g}, U(\mathfrak{l})) = \prod_{k \in \mathbb{Z}} H_{k+n}(\mathfrak{n}, \mathbb{C})$

Proof. The complex $ToT^{per}(\Omega^{\mathfrak{n}^*}(\mathfrak{g}^*))$ is a direct product of complexes 88 ■

$C^\bullet(\mathfrak{n}, \mathbb{C})$ is a subalgebra of $C^\bullet(\mathfrak{g}, \mathbb{C})$. We have an action of $C^i(\mathfrak{n}, \mathbb{C})$ on $C_n(\mathfrak{g}, U(\mathfrak{l}))$. The action is compatible with differential d . Moreover it commutes with d_{dR} .

From this we conclude that $C^\bullet(\mathfrak{n}, \mathbb{C})$ acts on the total complex 88.

The following bicomplex is a specialization of 88 ($\mathfrak{g} = L, \mathfrak{l} = TYM$).

$$\begin{array}{ccccc}
 \cdots & & \cdots & & \cdots \\
 \downarrow d & & \downarrow d & & \downarrow d \\
 \text{Sym}^0(TYM) \otimes \Lambda^2(L) & \xleftarrow{d_{dR}^L} & \text{Sym}^1(TYM) \otimes \Lambda^1(L) & \xleftarrow{d_{dR}^L} & \text{Sym}^2(TYM) \otimes \Lambda^0(L) \\
 \downarrow d & & \downarrow d & & \\
 \text{Sym}^0(TYM) \otimes \Lambda^1(L) & \xleftarrow{d_{dR}^L} & \text{Sym}^1(TYM) \otimes \Lambda^0(L) & & \\
 \downarrow d & & & & \\
 \text{Sym}^0(TYM) \otimes \Lambda^0(L) & & & &
 \end{array} \tag{89}$$

It leads to a spectral sequence

$$H_i(L, \text{Sym}^j(TYM)) \Rightarrow H_{i+2j}(\mathfrak{su}(\mathfrak{h}), \mathbb{C}) \tag{90}$$

Observe that the spaces $\text{Sym}^j(TYM) \otimes \Lambda^i(L)$ for fixed j form columns of 89. It makes sense to talk about cocycles of $C_i(L, \text{Sym}^j(TYM))$ as of elements of 89.

Proposition 57 *Every element in the image of the maps $\iota_s, s = 1 \dots 5$ in (54) for $i \geq 4$ can be modified to a cocycle of the total complex 89.*

Proof. The maps d and d_{dR} are $Spin(10)$ -equivariant. We can utilize this to solve “tic-tac” process: pick an element $a = a_0$ whose class is from the image of ι_s , moreover it belongs to an irreducible $Spin(10)$ subrepresentation of $\text{Sym}^j(TYM) \otimes \Lambda^i(L)$. The representation must be one of in (54). The element $d_{dR}a$ belongs an irreducible representation of the same isotopic type. From proposition (22) we know that it must be homologous to zero. We find a_1 again in irreducible subrepresentation of $\text{Sym}^{j-1}(TYM) \otimes \Lambda^{i+1}(L)$ such that $da_1 = d_{dR}a$. We continue this process until we end up in $\text{Sym}^0(TYM) \otimes \Lambda^{i+j}(L)$. Let $a_0 = a$. The sequence of elements (a_0, a_1, \dots, a_j) is a cocycle of 89. ■

Lemma 58 *The maps $\iota_j, 1 \leq j \leq 5$ in (54) are embeddings for $i \geq 0$.*

Proof. It is given in section (7.2). ■

Denote the image of ι_j in degree i by A_j^i

Lemma 59 *There is a cocycle representing any class in $\text{Im} \iota_j$ $i \geq 0$ which can be lifted to nontrivial cocycles of 89.*

Proof. The proof is a combination of proofs of proposition (57) and lemma (58).

Indeed, as in a proof of proposition (57), lift a cocycle a to a cocycle $a = a_0, \dots, a_i$. Apply multiplication on λ^{α_s} several times, until $\lambda^{\alpha_1} \dots \lambda^{\alpha_i} a_0$ lands in zero chains $\text{Sym}^j(TYM)$. As we know from lemma (58), this element projects nontrivially into $H_0(L, \text{Sym}^j(TYM))$. Thus the cochain $\lambda^{\alpha_1} \dots \lambda^{\alpha_i}(a_0, \dots, a_i)$ has a nontrivial class in cohomology. From this we conclude that so does the class (a_0, \dots, a_i) .

■

The cohomology of the total complex 89 is easy to compute.

Remark 60 *According to proposition (55) applied to $\mathfrak{g} = L, \mathfrak{l} = TYM$, the cohomology of 89 is equal to $H_\bullet(\mathfrak{su}\mathfrak{sh}, \mathbb{C})$.*

According to lemma (59) every class in $\text{Im} \iota_j$ gives a nontrivial contribution to $H_\bullet(\mathfrak{su}\mathfrak{sh}, \mathbb{C})$. We see that the image of the maps ι_j cover almost all cohomology of 89 (see the table (83) in appendix).

The exceptional representations in homology $H_\bullet(\mathfrak{su}\mathfrak{sh}, \mathbb{C})$ not covered by the above construction are

$$\begin{aligned} [0, 0, 0, 2, 0] &\subset H_9(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), [0, 0, 0, 1, 0] \subset H_8(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), \\ [0, 0, 0, 0, 0] &\subset H_7(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), [1, 0, 0, 1, 0] \subset H_7(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), \\ [1, 0, 0, 0, 0] &\subset H_6(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), [0, 1, 0, 0, 0] \subset H_5(\mathfrak{su}\mathfrak{sh}, \mathbb{C}), \\ [0, 0, 0, 0, 0] &\subset H_3(\mathfrak{su}\mathfrak{sh}, \mathbb{C}) \end{aligned} \tag{91}$$

To do this we need to prove the following

Proposition 61 $H_3(L, \text{Sym}^j(TYM)) = A_j^3$, $H_2(L, \text{Sym}^j(TYM)) = A_j^2$

Proof. The proof is given in section (7.6). ■

Proposition 62 *The spectral sequence (90) collapses in E_2 -term.*

Proof. We already know that classes in A_j^i survive to E_∞ . Due to proposition (61) and dimension reasons the spectral sequence collapses. ■

Corollary 63 *The differential*

$$d_{dR} : H_1(L, \text{Sym}^j(TYM)) \rightarrow H_2(L, \text{Sym}^{j-1}(TYM)) \quad (92)$$

is zero for $j \geq 0$.

Proof. According to proposition (61) $H_2(L, \text{Sym}^j(TYM)) = A_j^2$. This groups survives to E_∞ , thus the map (92) is trivial.

■

Corollary 64 *There are linear subspaces*

$$\begin{aligned} [0, 0, 0, 1, 0] &\subset H_{0,13}(L, \text{Sym}^4(TYM)) \\ [1, 0, 0, 0, 0] &\subset H_{0,10}(L, \text{Sym}^3(TYM)) \\ \mathbb{C}_{1,4} &\subset H_{1,4}(L, TYM) \\ [0, 1, 0, 0, 0] &\subset H_{1,8}(L, \text{Sym}^2(TYM)) \\ \mathbb{C}_{1,12} &\subset H_{1,12}(L, \text{Sym}^3(TYM)) \\ [1, 0, 0, 1, 0] &\subset H_{1,11}(L, \text{Sym}^3(TYM)) \\ [0, 0, 0, 2, 0] &\subset H_{1,14}(L, \text{Sym}^4(TYM)) \end{aligned} \quad (93)$$

$$\text{The space } \mathbb{C}_{1,4} \text{ is spanned by } \gamma_{1,4} = \sum_{\alpha} \lambda_{\alpha}^* \chi^{\alpha} \quad (94)$$

$$\text{the space } \mathbb{C}_{1,12} \text{ is spanned by } \gamma_{1,12} = \Gamma_{\alpha}^{[\beta[i_1, \dots, i_4]]} \lambda_{\alpha}^* \otimes \chi^{\alpha} \circ F_{i_1 i_2} \circ F_{i_3 i_4}$$

Where $F_{ij} = [v_i, v_j]$ and \circ is a graded symmetric product.

The spaces (93) belong to the kernel of d_{dR} . Denote a direct sum of the above subspaces of $H_i(L, \text{Sym}^j(TYM))$ by B_j^i

Fix n . The direct sum $B_j^{n-j} + A_j^{n-j}$ isomorphically project onto j -th cohomology of the complex $(H_{n-j}(L, \text{Sym}^j(TYM)), d_{dR})$.

Proof. Simple use of propositions (62, 61, 50) ■

Proposition 65 *The maps*

$$\pi_j^i H^{j+i, 2j-i}(L, \text{Sym}^j(TYM)) \rightarrow H^0(\mathcal{Q}, \Lambda^j(\mathcal{W})(j+i)) \quad (95)$$

$i, j \geq 0$ defined in (20) are surjections.

Proof. It is easy to see that the map π_j^i maps cocycle $\Gamma_{\alpha\beta}^s \lambda^\alpha \chi^\beta$ to a generator $v_s \in H^0(\mathcal{Q}, \Lambda^1(\mathcal{W})(1))$ $s = 1 \dots, 10$. The products of v_s and $\lambda^\alpha \in H^0(\mathcal{Q}, \mathcal{O}(1))$ generate $H^0(\mathcal{Q}, \Lambda^j(\mathcal{W})(j+i))$. We interpret λ^α as elements of $H^1(L, \mathbb{C})$.

We prove proposition using that π is an algebra homomorphism, the target has no zero divisors as an algebra over S and π_j^i is an isomorphism for $i \geq 4$ ■

Proposition 66 *The following maps defined in proposition (20) are surjections.*

$$\begin{aligned} H^{3,12}(L, \text{Sym}^2(TYM)) &\rightarrow H^9(\mathcal{Q}, \Lambda^2(\mathcal{W})(-6)) = \mathbb{C} \\ H^{2,8}(L, \text{Sym}^3(TYM)) &\rightarrow H^1(\mathcal{Q}, \Lambda^3(\mathcal{W})(1)) = \mathbb{C} \\ H^{3,16}(L, \text{Sym}^3(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^3(\mathcal{W})(-7)) = [0, 1, 0, 0, 0] \\ H^{4,18}(L, \text{Sym}^4(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^4(\mathcal{W})(-6)) = [1, 0, 0, 0, 0] \\ H^{3,19}(L, \text{Sym}^4(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^4(\mathcal{W})(-7)) = [1, 0, 0, 1, 0] \\ H^{5,20}(L, \text{Sym}^4(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^5(\mathcal{W})(-5)) = \mathbb{C} \\ H^{4,21}(L, \text{Sym}^5(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^5(\mathcal{W})(-6)) = [0, 0, 0, 1, 0] \\ H^{3,22}(L, \text{Sym}^5(TYM)) &\rightarrow H^{10}(\mathcal{Q}, \Lambda^5(\mathcal{W})(-7)) = [0, 0, 0, 2, 0] \end{aligned} \quad (96)$$

Proof. In long exact sequence (51) all linear spaces $H^i(\mathcal{Q}, \Lambda^j(\mathcal{W})(k))$ from (96) are mapped to a zero space. ■

Consider a complex 88 for a pair $TYM \subset YM$.

Proposition 67 *The spectral sequence of bicomplex 88 for a pair $TYM \subset YM$ has the only nontrivial differential in E_2 -term, which maps a class of Lagrangian in $H_0(YM, \text{Sym}^2(TYM))$ to the generator of $H_3(YM, \mathbb{C})$*

Proof. We leave the proof of this proposition as an exercise for the reader. Most of the differentials are zero because of vanishing of $H_3(YM, \text{Sym}^k(TYM))$. ■

Proposition 68 *The spectral sequence of proposition (67) splits into exact sequences:*

$$\begin{aligned}
0 \rightarrow \Lambda^{2i}(V) \rightarrow H_0(YM, \text{Sym}^i(TYM)) &\xrightarrow{d_{dR}^{YM}} H_1(YM, \text{Sym}^{i-1}(TYM)) \xrightarrow{\nu} \\
&\xrightarrow{\nu} \Lambda^{2i-1}(V) \rightarrow 0 \quad i > 3 \\
0 \rightarrow \Lambda^6(V) \rightarrow H_0(YM, \text{Sym}^3(TYM)) &\xrightarrow{d_{dR}^{YM}} H_1(YM, \text{Sym}^2(TYM)) \xrightarrow{\nu} \\
&\xrightarrow{\nu} \Lambda^5(V) + H_2(YM, TYM) \rightarrow 0 \\
0 \rightarrow \Lambda^4(V) \rightarrow H_0(YM, \text{Sym}^2(TYM))/\mathbb{C} &\xrightarrow{d_{dR}^{YM}} H_1(YM, \text{Sym}^1(TYM)) \xrightarrow{\nu} \quad (97) \\
&\xrightarrow{\nu} \Lambda^3(V) + H_2(YM, \mathbb{C}) \rightarrow 0 \\
0 \rightarrow \Lambda^2(V) \rightarrow H_0(YM, TYM) &\xrightarrow{d_{dR}^{YM}} H_1(YM, \mathbb{C}) \xrightarrow{\nu} \\
&\xrightarrow{\nu} \Lambda^1(V) \rightarrow 0
\end{aligned}$$

Remark 69 *The map d_{dR}^{YM} corresponds to taking equations of motion of a Lagrangian. The later is an element of zero homology group. The kernel of the map d_{dR}^{YM} are topological term of the Lagrangian.*

Remark 70 *We can say that the elements of $H_2(YM, \text{Sym}(TYM))$ which do not map to zero under ν are nonlagrangian deformations of equations of motion.*

Remark 71 *By proposition (76) $\text{res}\delta d_{dR}^L f$, $f \in H_0(L, \text{Sym}(TYM))$ automatically belongs to the image of d_{dR}^{YM} and hence is a Lagrangian deformation. The cocycle $c_{2,8} \in H^{2,8}(L, \text{Sym}^3(TYM))$ has Poincare dual in $H_{1,16}(YM, \text{Sym}(TYM))$. The later element has image under ν equal to zero in (97) due to degree considerations. The same arguments applies to cocycles $\delta\gamma_{1,4}, \delta\gamma_{1,12}$.*

7.5 On the structure of the connecting differential δ from (77)

The map δ plays an important role in our construction of infinitesimal deformations. In this section we make a preliminary study of δ and identify it with a cocycle in $HH^\bullet(S, S \otimes S)$. Let N be an L -module. The table below represents E_1 term of a spectral sequence of a bicomplex $E_1^{i,j} \Rightarrow H^{i+j}(\Omega\mathcal{N}(l))$. See definition (38) for explanation of notations.

		N_{-l-10}		N_{-l-9}															
10	...	\otimes	\rightarrow	\otimes	\rightarrow	N_{-l-8}	...	0		0				0		...			
		S_2^*		S_1^*															
9	...	0		0		0	...	0		0				0		...			
...			
1	...	0		0		0	...	0		0				0		...			
0	...	0		0		0	...	N_{-l}	\rightarrow	\otimes	\rightarrow	\otimes		...					
										S_1		S_2							
		$-l - 10$		$-l - 9$		$-l - 8$		$-l$		$-l + 1$		$-l + 2$							

Let us describe the connecting differential of (7.5) in analytic terms. Fix $Spin(10, \mathbb{R})$ -invariant Kahler metric on \mathcal{Q} and on $\mathcal{O}(1)$. Let $p_k : \Omega^{0,k}(l) \rightarrow H^k(\mathcal{Q}, \mathcal{O}(l))$ be the orthogonal projection onto cohomology. We have $p_k = 0$ for $k \neq 0, 10$. Let us set $p = \bigoplus_{k=0}^{10} p_k$. One can choose $SO(10)$ -equivariant homotopy $L : \Omega^{0,k}(l) \rightarrow \Omega^{0,k-1}(l)$ which satisfies the following set of properties: $L^2 = 0$, $\{\bar{\partial}, L\} = Id - p$, where Id is the identity transformation.

Suppose $a = \sum a_i \otimes \omega_i \in N \otimes S^*$ is a representative of a cohomology class of one of cohomology groups in the tenth row of table (7.5). Let $e = \sum_{\alpha} \theta_{\alpha} \lambda^{\alpha}$. The element $b_1 = \sum \theta_{\alpha} a_i \otimes \lambda^{\alpha} \omega_i$ is $\bar{\partial}$ -coboundary. Define an element $c_1 = \sum \theta_{\alpha} a_i \otimes L(\lambda^{\alpha} \omega_i)$. By construction $\bar{\partial} c = b_1$. Iterating this construction, we get an element

$$\theta_{\alpha_1} \theta_{\alpha_2} \dots \theta_{\alpha_{11}} a_i \dots \otimes \lambda^{\alpha_1} L(\lambda^{\alpha_2} \dots L(\lambda^{\alpha_{11}} \omega_i) \dots)) \quad (98)$$

, which is δ differential of element a .

There is an algebraic description of the differential δ .

To give it we need to make a digression. Let A be an algebra N', N'' are two modules. We can define groups $Ext_A^i(N', N''), i > 0$ following Yoneda. Consider an acyclic complex of modules

$$0 \rightarrow N'' \rightarrow P_1 \rightarrow \dots \rightarrow P_i \rightarrow N' \rightarrow 0$$

These exact sequences form a semigroup. There is an obvious notion of an equivalence of such sequences and a notion of a "trivial" sequences. After factorization with respect to equivalence relation and after killing all trivial elements we get groups $Ext_A^i(N', N'')$ (see [12] for details).

We utilize the Dolbeault complex $\bigoplus_{l \in \mathbb{Z}} \Omega^{0, \bullet}(l)$ to construct an element in $Ext_S^{11,8}(S^*, S)$. The Dolbeault differential $\bar{\partial}$ is linear with respect to multiplication on elements of $S \subset \bigoplus_{l \in \mathbb{Z}} O(l) \otimes \Omega^{0,0}$. It means that we can interpret a complex

$$0 \rightarrow S \rightarrow \bigoplus_{l \in \mathbb{Z}} \Omega^{0,0}(l) \rightarrow \dots \rightarrow \bigoplus_{l \in \mathbb{Z}} \Omega^{0,10}(l) \rightarrow S^* \rightarrow 0 \quad (99)$$

as an element of $Ext_S^{11,8}(S^*, S)$.

According to [4] for any algebra A and two left modules M, N there is an isomorphism $Ext_A^n(M, N) = HH^n(A, N \otimes M^*)$. We conclude $Ext_S^n(S^*, S) = HH^n(S, S \otimes S)$. The later group can be computed via Koszul resolution:

Proposition 72 *The cohomology $HH^n(S, S \otimes S)$ can be computed as cohomology of the complex $U(L) \otimes S \otimes S$. The differential is defined for homogeneous elements by the formula*

$$d(a \otimes b \otimes c) = (\theta_\alpha a \otimes \lambda^\alpha b \otimes c - (-1)^{\bar{a}} a \theta_\alpha \otimes b \otimes \lambda^\alpha c) \quad (100)$$

Proof. is similar to (18) and (19) ■

Proposition 73 *The cohomology $HH^n(S, S \otimes S)$ is equal to 0 for $n \neq 11$ and S for $n = 11$*

Proof. We need to start with a remark that this statement is true for coordinate ring of any smooth affine variety. In non smooth case this statement is typically not correct, which makes this proposition a bit surprising.

We do computations with $U(L) \otimes S \otimes S$ as it was explained in (72). The statement of proposition easily follows from the following fact:

Lemma 74 *The cohomology of the complex $U(L) \otimes S$ with differential equal to left multiplication on $\lambda^\alpha \theta_\alpha$ is equal to \mathbb{C} . The nontrivial cocycle has bidegree $(11, 3)$.*

Proof. As usual, we use localization arguments. Fix $\theta \in L_1$, such that $\theta^2 = 0$. The cohomology of $U(L)$ with differential-on θ are trivial (use a spectral sequence associated with PBW filtration to prove this statement).

The spectral sequence of the hypercohomology of the localized complex $(\mathcal{U}(\mathcal{L}), d)$ for $(U(L) \otimes S, d)$ converges to zero and collapses in E_1 . The only nontrivial higher differential $d_1 0$ in $E_3 = E_1 0$ establishes isomorphism between cohomology $U(L) \otimes S^*$ and $U(L) \otimes S$

The algebra S is Koszul. Thus by definition the complex $U(L) \otimes S^*$ has the only nontrivial cohomology in bidegree $(0, 0)$. The image of this cohomology class in $U(L) \otimes S$ is the one with bidegree $(11, 3)$. ■

■

Let N be an L -module. There is a sequence of maps

$$H_0(YM, N) \xrightarrow{i} H_0(L, N) \xrightarrow{\delta} H^3(L, N) \xrightarrow{res} H^3(YM, N) \xrightarrow{P} H_0(YM, N) \quad (101)$$

The map i is a map of zero homology of subalgebra on zero homology of algebra, δ is the differential (98), res is the restriction map from cohomology of algebra to cohomology of subalgebra, P is the Poincare isomorphism for the algebra YM .

Denote the composed map by ψ_0 .

Proposition 75 *Denote $x = \epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}}$ an element in $U(L)$. The map ψ_0 is defined by the formula*

$$\psi_0(n) = xm \quad (102)$$

Proof. If we change the map xm by adding to x an element \tilde{x} of the same degree but lower order in PBW filtration the map $xm + \tilde{x}m$ is still equal to xm .

Denote $YM(N) = \{n \in N | n = \sum_i l_i n_i, l_i \in YM, n_i \in N\}$. It is easy to see that on the level of elements of N , $\tilde{x}n \in YM(N)$, but by definition $H_0(YM, N) = N/YM(N)$

The map ψ_0 is completely determined by cocycle (99), which we interpret as an element of $H^3(L, U(L))$. The Lie algebra L acts on $U(L)$ by left multiplication.

We will reinterpret this cocycle through spectral sequence of extension $YM \subset L$. The E_1^{ij} -term is equal to $C^i(L/YM, H^j(YM, U(L)))$. It is easy to see that $U(L)$ is free as a $U(YM)$ -module under left multiplication (in fact this is true for any Lie algebra and subalgebra). We conclude that $H^j(YM, U(L)) = 0$, $j \neq 3$ and $\mathbb{C} \otimes_{U(YM)} U(L) = \Lambda(L_1)$ for $j = 3$. The $E_2^{i,3}$ is equal to cohomology of Koszul complex $H^\bullet(L_1, \Lambda(L_1)) = H^\bullet(\text{Sym}(L_1^*) \otimes \Lambda(L_1))$. The cohomology is one-dimensional and are represented by cocycle $\epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}} \in \Lambda^{16}(L_1)$. A lift of this class to a cocycle in $C^3(YM, U(L))$ and identification with $C_0(YM, U(L))$ involves ambiguities (chooses in lower order of PBW-filtration), which are do not affect the final map as we showed above. ■

There is a minor generalization of the map ψ . There are compositions of maps

$$H_i(YM, N) \xrightarrow{\iota} H_i(L, N) \xrightarrow{\delta} H^{3-i}(L, N) \xrightarrow{res} H^{3-i}(YM, N) \xrightarrow{P} H_i(YM, N) \quad (103)$$

We denote them by ψ_i .

Proposition 76 *The maps $\psi_i H_i(YM, \text{Sym}(TYM)) \rightarrow H_i(YM, \text{Sym}(TYM))$ commute with differential d_{dR}^{YM} .*

Proof. Any derivation of YM which preserves TYM acts on $H_i(YM, \text{Sym}(TYM))$ due to functionality. Due to the same functionality the action of derivations is compatible with d_{dR}^{YM} . It implies that d_{dR}^{YM} is compatible with the action of $x = \epsilon^{\alpha_1 \dots \alpha_{16}} \theta_{\alpha_1} \dots \theta_{\alpha_{16}}$. ■

7.6 Computation of $H^0(L, \text{Sym}^j(TYM)), H^1(L, \text{Sym}^j(TYM))$

According to long exact sequence (51) and proposition (58), we have an isomorphism $H_3(L, \text{Sym}^j(TYM))/A_j^3 = H^3(L, \text{Sym}^j(TYM))$. $H_2(L, \text{Sym}^j(TYM))/A_j^3 = H^1(L, \text{Sym}^j(TYM))$.

The present section content is a proof of the following lemma

Lemma 77 $H^1(L, \text{Sym}^j(TYM)) = 0, j \geq 2$ and $H^1(L, TYM) = [1, 0, 0, 0, 0]$ and $H^0(L, \text{Sym}^j(TYM)) = 0, j \geq 1$, elements $H^1(L, TYM)$ have degree two.

For the proof we adopt a method developed in [14] where we treated a similar problem in a more simple context of pure Yang-Mills theory.

There is a short exact sequence of algebras where L_1 is an abelian Lie algebra in degree one.

For an estimate of $H^i(L, \text{Sym}^k(TYM))$ we can use a Serre-Hochschild spectral sequence (62). If we manage to prove that $H^1(YM, \text{Sym}^k(TYM)) = 0, k \geq 2$, $H^1(YM, TYM) = [1, 0, 0, 0, 0] + [1, 0, 0, 1, 0]$, $H^0(YM, \text{Sym}^k(TYM)) = 0, k \geq 1$ we shall be done.

The text below follows very close to [14]. We omit the proofs which are close to their bosonic counterparts.

We proceed as in [14] replacing $\text{Sym}(TYM)$ by $U(TYM)$. The universal enveloping of TYM is a free algebra generated by \mathbb{Z}_2 graded linear space M . Define a filtration by powers of augmentation ideal $I \subset U(TYM)$. The adjoint action of YM on $\bigoplus_{i \geq 0} Gr_i I^i / I^{i-1} = \bigoplus_{i \geq 0} M^{\otimes i}$ factors through abelianization $Ab(YM) = L_2 + L_3$. The component L_3 of $Ab(YM)$ acts on $\bigoplus_{i \geq 0} M^{\otimes i}$ trivially.

Let us describe the linear space M , which is automatically $U(L_2) = \text{Sym}(V)$ -module. It consists of two components M_0 and M_1 (see [16]). They admit a geometric interpretation.

Consider a nonsingular quadric $\mathcal{X} \subset \mathbf{P}^9$ defined by equation $q = 0$ where

$$q = x_i^2 \in \text{Sym}(V) \quad (104)$$

The polarization of defining equation is the bilinear form, used in definition of algebra YM . Let T be the tangent bundle to X . Define $M_0 = \bigoplus_{i \geq 0} H^0(\mathcal{X}, T(i))$.

The quadric is a homogeneous space of $Spin(10)$. The Levi subgroup of the stabilizer of a point $x \in \mathcal{X}$ is $\mathbb{C}^* \times Spin(8)$. The group $Spin(8)$ acts on the fiber via fundamental 8-dimensional representation.

There is a famous triality for the group $Spin(8)$. It collects all representations of $Spin(8)$ in orbits S_3 action.

The orbit of the defining representation v of $Spin(8)$ contains s^+ and s^- -the spinor representations. They also have dimension eight.

To define M_1 we take s^+ and induce a holomorphic vector bundle on \mathcal{X} , which we denote by f^+ . Then $M_1 = \bigoplus_{i \geq 0} H^0(X, f^+(i))$.

Fix a YM -module N . The following is a complex constructed from a free $U(YM)$ -resolution of trivial YM -module \mathbb{C} . The complex computes $H^i(YM, N) = H_{3-i}(YM, N)$

$$N \xrightarrow{d_1} N \otimes (V + S^*) \xrightarrow{d_2} N \otimes (V + S) \xrightarrow{d_3} N \quad (105)$$

$$\begin{aligned} d_1(f \otimes c) &= v_i f \otimes v_i^* + (-1)^{\tilde{f}} \chi^\alpha f \otimes \chi_\alpha^* \\ d_2(f \otimes v_k^*) &= (v_i v_i f \otimes v_k + v_i v_k f \otimes v_i - 2v_k v_i f \otimes v_i) + \\ &\quad + (-1)^{\tilde{f}} \Gamma_{\alpha\beta}^k \chi^\alpha f \otimes \chi^\beta \\ d_2(h \otimes \chi_\alpha^*) &= -\Gamma_{\alpha\beta}^i \left(v_i h \otimes \chi^\beta - (-1)^{\tilde{h}} \chi^\beta h \otimes v_i \right) \\ d_3(f \otimes v_i + h \otimes \chi_\alpha) &= v_i f + (-1)^{\tilde{h}} \chi_\alpha h \end{aligned}$$

In case of module $M^{\otimes n}$ the complex (105) splits into sum of two complexes:

$$\langle c^* \rangle \otimes M^{\otimes n} \xrightarrow{d_1^{YM}} V \otimes M^{\otimes n} \xrightarrow{d_2^{YM}} V \otimes M^{\otimes n} \xrightarrow{d_3^{YM}} \langle c \rangle \otimes M^{\otimes n} \quad (106)$$

$$S^* \otimes M^{\otimes n} \xrightarrow{d_2^D} S \otimes M^{\otimes n} \quad (107)$$

since the action of odd generators of YM on M is trivial. We denote the above direct sum by $SC(M^{\otimes n})$

Let N be a $\text{Sym}(V)$ -module. Consider a complex $N \otimes \Lambda(V)$. The linear space $V \subset \Lambda(V)$ has a basis $\varsigma_1, \dots, \varsigma_{10}$ in degree one. A linear subspace $V \subset \text{Sym}(V)$ has a basis x_1, \dots, x_{10} in degree two. The differential d in the complex $N \otimes \Lambda(V)$

is defined by the formula $d(n \otimes \varsigma_i) = x_i n$ on generators and extended by the Leibniz rule.

Proposition 78 *There is a long exact sequence*

$$\begin{aligned} 0 \rightarrow H_3(V, M^{\otimes j}) &\xrightarrow{S_j} H_1(V, M^{\otimes(j+1)}) \xrightarrow{B_j} H_2(YM, M^{\otimes j}) \xrightarrow{I_j} \\ &\rightarrow H_2(V, M^{\otimes j}) \xrightarrow{S_j} H_0(V, M^{\otimes(j+1)}) \xrightarrow{B_j} H_1(YM, M^{\otimes j}) \xrightarrow{I_j} H_1(V, M^{\otimes j}) \rightarrow 0 \end{aligned} \quad (108)$$

and isomorphisms

$$H_3(YM, M^{\otimes j}) = 0 \quad j \geq 1 \quad (109)$$

$$H_0(YM, M^{\otimes j}) = H_0(V, M^{\otimes j}) \quad (110)$$

$$H_s(V, M^{\otimes j}) = \Lambda^{2j+s}[V] \quad s \geq 2, \text{ except } s = 2, j = 1 \quad (111)$$

For $j = 1$ we have

$$\begin{aligned} 0 \rightarrow \Lambda^4 V &\rightarrow H_2(V, M_0) \rightarrow \mathbb{C} \rightarrow 0 \\ 0 \rightarrow \Lambda^3 V &\rightarrow H_1(V, M_0) \rightarrow V \rightarrow 0 \quad H_1(V, M_1) = S \\ \Lambda^2 V &= H_0(V, M_0) \quad H_1(V, M_1) = S^* \end{aligned} \quad (112)$$

There is $\text{Sym}(V)$ -linear map

$$\delta^c : H_i(YM, M^{\otimes j}) \rightarrow H_{i+1}(YM, M^{\otimes(j-1)}) \quad (113)$$

Denote composition $B_{j-1} \circ I_j = \delta_j^c$. Then $\delta_{j-1}^c \circ \delta_j^c = 0$

Proof. The proof repeats the proof of proposition 10 in [14]. To compute $H_i(V, M_1)$ we used a free $\text{Sym}(V)$ -resolution of length two from [16]:

$$M_1 \leftarrow S \otimes \text{Sym}(V) \leftarrow S^* \otimes \text{Sym}(V) \leftarrow 0 \quad (114)$$

Suppose s_α is a basis S and s^α dual basis of S^* . Then $ds^\alpha = \Gamma_i^{\alpha\beta} x^i s_\alpha$.

■

Corollary 79 $H_i(V, M^{\otimes n}) = 0$ if $n \geq 2$ and $i = 9, 10$.

Proposition 80 For $n \geq 2$ the operator of multiplication on q in $M^{\otimes n}$ has no kernel.

Proof. The same as a proof of proposition 17 in [14]. ■

Proposition 81 For $n \geq 2$ the first cohomology of the complex $SC(M^{\otimes n})$ equal to $H^1(YM, M^{\otimes n})$ are trivial.

Proof. The complex $SC(M^{\otimes n})$ splits into sum of two: the complex 106 $C(M^{\otimes n})$ (defined in [14] equation 17) and the complex 107.

The first cohomology of the complex $C(M^{\otimes n})$ is equal to the group $H(M^{\otimes n})$ (see definition 13 [14]). Since $H_i(V, M^{\otimes n}) = 0$, $i = 9, 10$ and operator of multiplication on q in $M^{\otimes n}$ has no kernel then by proposition 15 [14] $H(M^{\otimes n}) = 0$

The complex 107 admits a homotopy $M^{\otimes n} \otimes S^* \xleftarrow{H} M^{\otimes n} \otimes S$, defined by the formula $H(m \otimes s_\alpha) = \Gamma_{\alpha\beta}^i x_i m \otimes s^\beta$. The composition Hd_D is proportional to an operator of multiplication on q , which has zero kernel in $M^{\otimes n}$, $n \geq 2$. It implies that d_D has no kernel.

Thus the complex $SC(M^{\otimes n})$ has vanishing first cohomology. ■

Proposition 82 The zero cohomology of $SC(M^{\otimes n})$, equal to $H^0(YM, M^{\otimes n})$ are zero for $n \geq 1$.

Proof. Zero cohomology of the complex is equal to $H_{10}(V, M^{\otimes n})$. The later group is zero by corollary 79. ■

Proposition 83 Denote

$$\begin{aligned} A &= \bigoplus_{i \geq 0} ([i, 2, 0, 0, 0] + [i + 2, 0, 0, 0, 0] + [i + 1, 0, 1, 0, 0]) \\ B &= \bigoplus_{i \geq 0} ([i, 1, 0, 0, 1] + [i + 1, 0, 0, 1, 0]) \\ C &= \bigoplus_{i \geq 0} ([i + 1, 0, 0, 0, 0] + [i, 0, 0, 0, 2] + [i, 0, 1, 0, 0]) \end{aligned} \tag{115}$$

Then

$$\begin{aligned}
H_0(V, M_0 \otimes M_0) &= A + \mathbb{C} + \Lambda^2(V) + \Lambda^4(V) \\
H_0(V, M_0 \otimes M_0) &= A + V + \Lambda^3(V) + \Lambda^5(V) \\
H_i(V, M_0 \otimes M_0) &= \Lambda^{i+4}(V) \quad i \geq 2 \\
H_0(V, M_0 \otimes M_1) &= B + [0, 0, 0, 0, 1] \quad H_1(V, M_0 \otimes M_1) = B + [0, 0, 0, 1, 0] \\
H_i(V, M_0 \otimes M_1) &= 0 \quad i \geq 2 \\
H_0(V, M_1 \otimes M_1) &= C \quad H_1(V, M_1 \otimes M_1) = C + \mathbb{C} \\
H_i(V, M_1 \otimes M_1) &= 0 \quad i \geq 2
\end{aligned} \tag{116}$$

Proposition 84 $H^0(YM, IU(TYM)) = 0$, $H^1(YM, IU(TYM)) = [1, 0, 0, 0, 0] + [0, 0, 0, 1, 0]$. The elements of $[1, 0, 0, 0, 0]$ have degree two, of $[0, 0, 0, 1, 0]$ degree one

Proof. The universal enveloping algebra $U(TYM) = T(M)$ admits a filtration by powers of the augmentation ideal $I \subset U(TYM)$. The adjoint action of YM preserves I , hence the filtration $F^i = I^{\times i}$. We plan to compute cohomology $H^i(YM, U(TYM))$ using a spectral sequence of mentioned filtration.

The E_2 term of it is equal to $E_2^{ij} = H^{i+j}(YM, M^{\otimes j})$. We computed $H^1(YM, M^{\otimes j})$ for $j \geq 1$. According to proposition (81) the groups are equal to zero for $j \geq 2$. We examine the differential in the spectral sequence on the group $H^1(YM, M)$. The differential δ acts:

$$\delta : H_2(YM, M) \rightarrow H_1(YM, \Lambda^2[M]) \subset H_1(YM, M^{\otimes 2}) \tag{117}$$

$H^0(YM, IU(TYM)) = 0$ due to proposition (82).

In the following part of the section we formulate essential lemmas needed for the proof that the kernel of δ is equal to $V + S^*$. ■

Proposition 85 The map $B_1 : H_1(V, M \otimes M) \rightarrow H_2(YM, M)$ is surjective.

Proposition 86 The map δ is an embedding on the image $\text{Im}(H_1(V, \Lambda^2(M)) \subset H_2(YM, M)$. The composition $\delta^c \circ \delta : H_2(YM, M) \rightarrow H_2(YM, M)$ is a projection on $\text{Im}(H_1(V, \Lambda^2(M)))$

Definition 87 For a (graded) Lie algebra \mathfrak{g} the group $D(\mathfrak{g})$ is a group of inner \mathfrak{g} -invariant co-products: $D(\mathfrak{g}) = \text{Sym}^2(\mathfrak{g})_{\mathfrak{g}}$. The linear space $D(\mathfrak{g})$ is generated by elements $a \circ b, a, b \in \mathfrak{g}$. Subject to relation $[a, b] \circ c + b \circ [a, c] = 0, a \circ b = b \circ a$ and the symbol is linear with respect to each of the arguments.

We would like to specialize construction of definition (87) to algebra $\mathfrak{h} = TYM/[TYM, [TYM, TYM]]$. The algebra \mathfrak{h} is a direct sum of two linear spaces $M + \Lambda^2(M)$. The linear space $\Lambda^2(M)$ is the center. The commutator $[\cdot, \cdot] : M \wedge M \rightarrow \Lambda^2(M)$ is an isomorphism.

Proposition 88 A linear space $D(\mathfrak{h})$ is a $\text{Sym}(V)$ -module. There is a short exact sequence of modules

$$0 \rightarrow \Lambda^3(M) \rightarrow D(\mathfrak{h}) \rightarrow \text{Sym}^2(M) \rightarrow 0 \quad (118)$$

Proposition 89 There is a commutative diagram

$$\begin{array}{ccc} H_2(YM, M) & \xrightarrow{\delta} & H_1(YM, M^{\otimes 2}) \\ \uparrow B_2 & & \uparrow B_3 \\ H_1(V, \text{Sym}^2(M)) & \xrightarrow{\tilde{\delta}} & H_0(V, \Lambda^3(M)) \end{array} \quad (119)$$

The map $\tilde{\delta}$ is the boundary map corresponding to extension $D(\mathfrak{h})$. The map B_3 in the diagram above has a trivial kernel. The kernel of map B_2 in (119) has kernel equal to $\Lambda^5(V)$.

Proposition 90 The kernel of $\tilde{\delta}$ is $\Lambda^5(V) + V + S^*$.

Proof. We send the reader to [14]. An important comment is that a nontrivial $\text{Spin}(10)$ -invariant cocycle $a \in H_1(V, M_1 \otimes M_1)$ is graded anti-symmetric and can be ignored in the present discussion. ■

Proposition 91 Elements of $[0, 0, 0, 1, 0] \subset H^{1,1}(YM, TYM)$ can not be lifted to $H^{1,1}(L, TYM)$

Proof. According to long exact sequence (51) the linear space $H^{1,1}(L, TYM)$ is equal to zero. ■

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